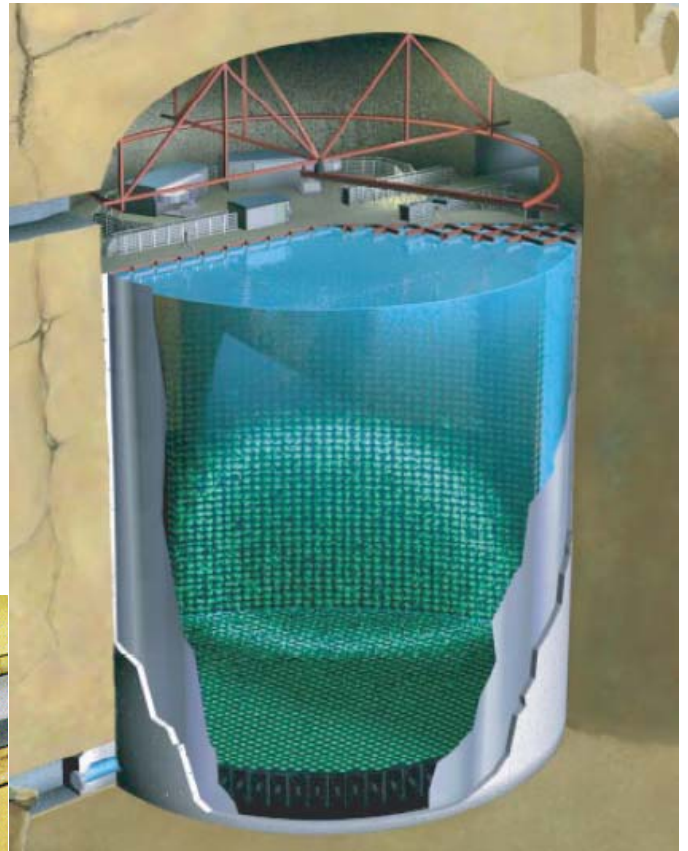
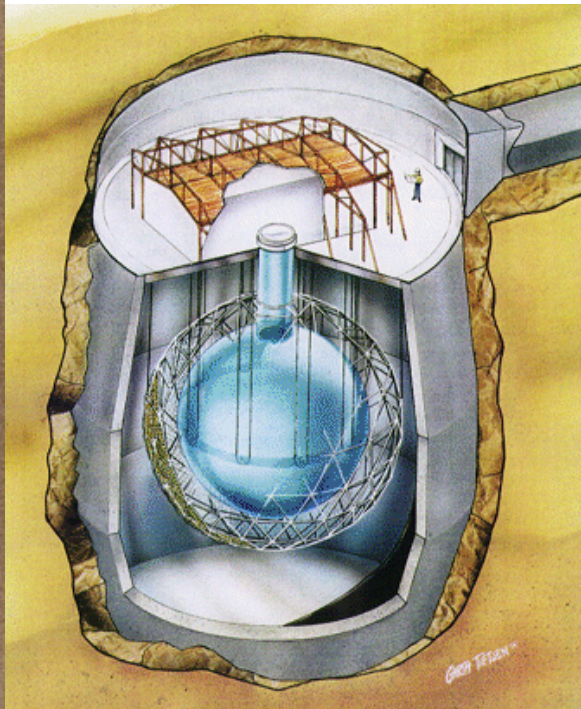
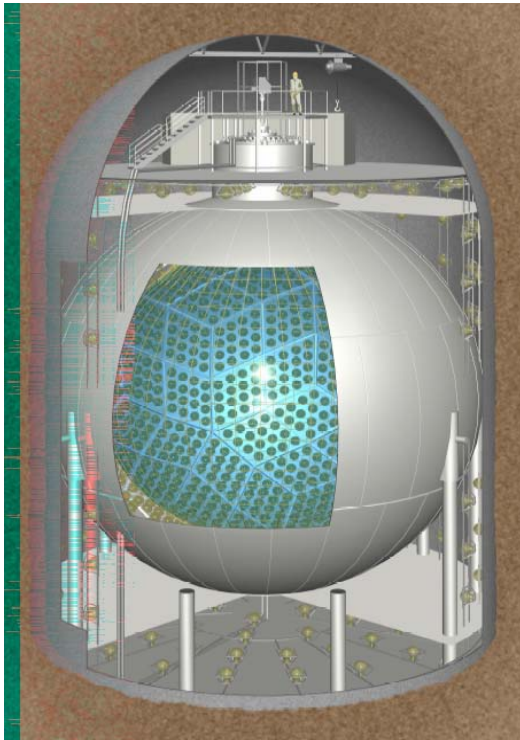


# Neutrinos; Looking forward to the future

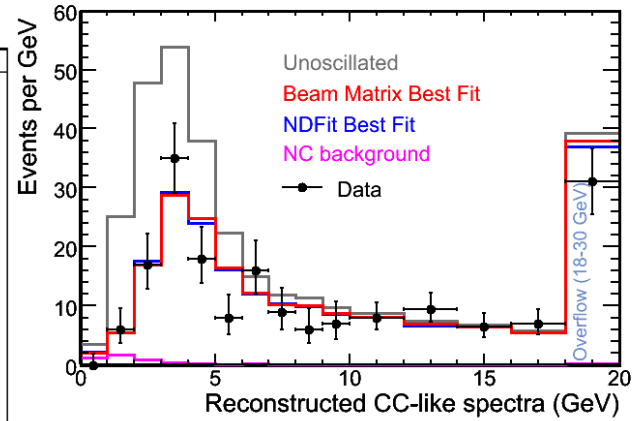
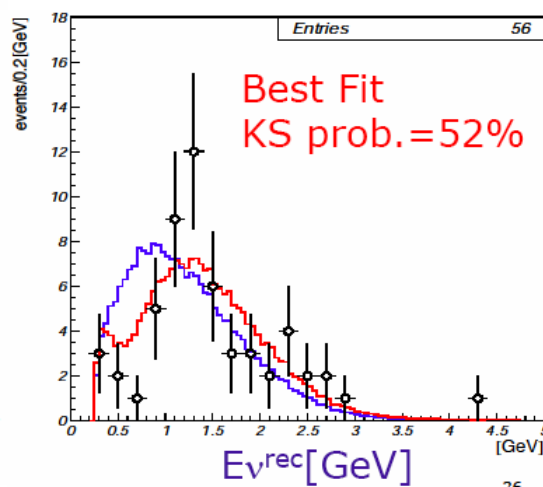
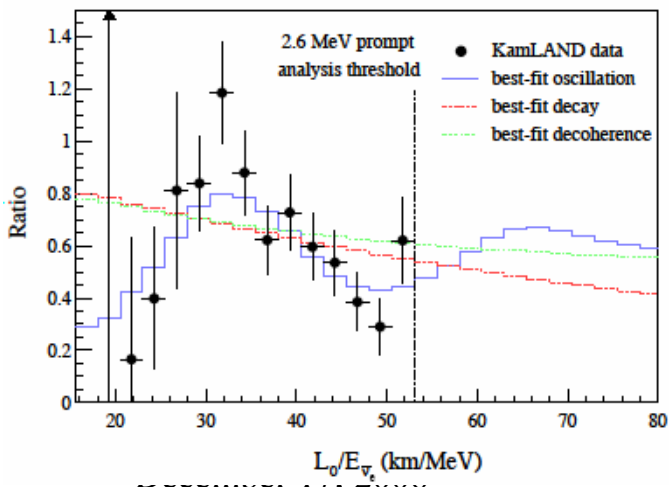
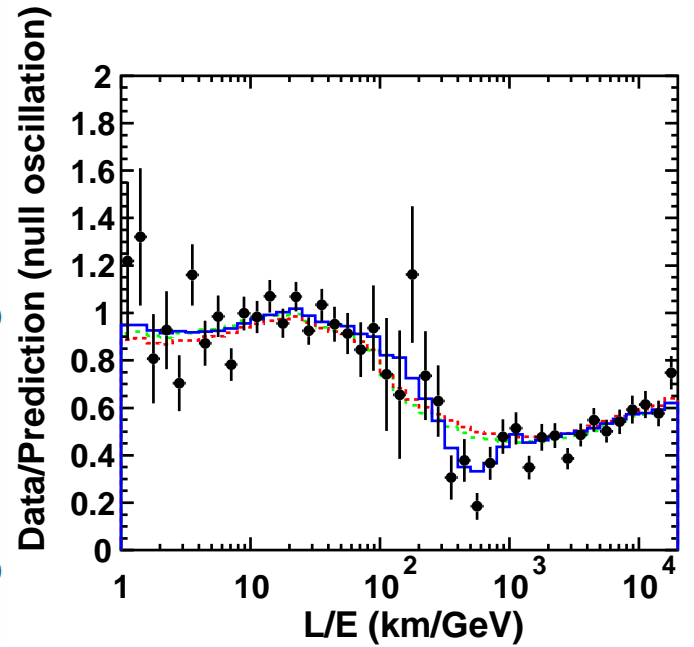
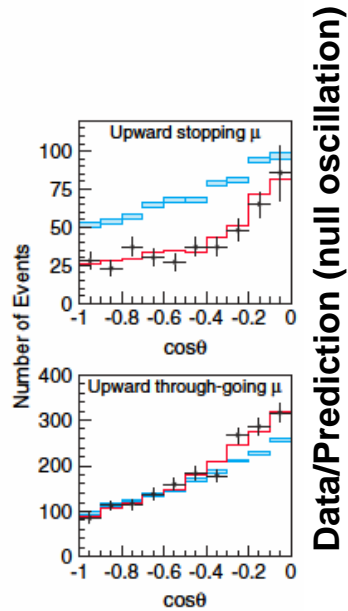
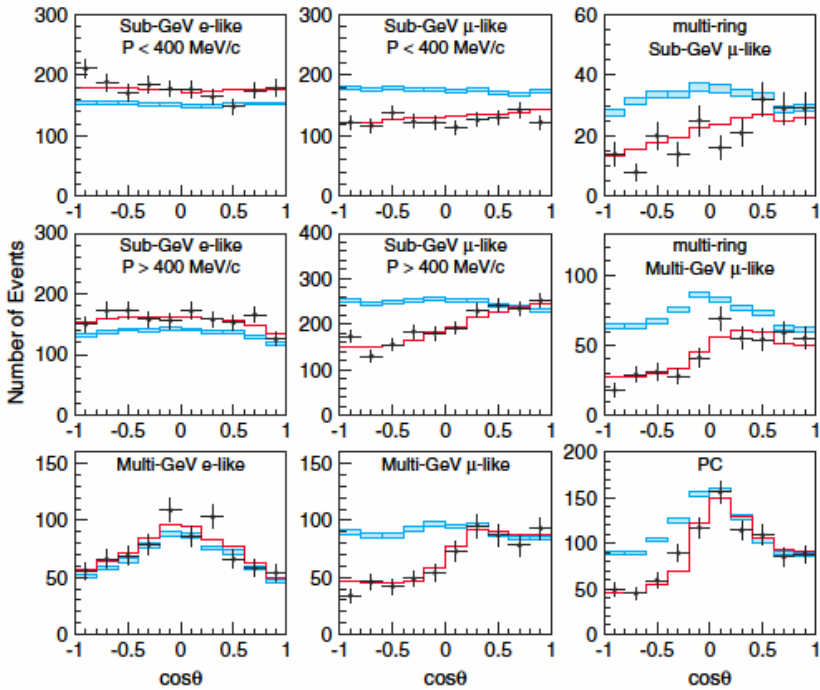


Hisakazu Minakata  
Tokyo Metropolitan University

In the last several years  
we have experienced the  
most exciting era in  
 $\nu$  physics

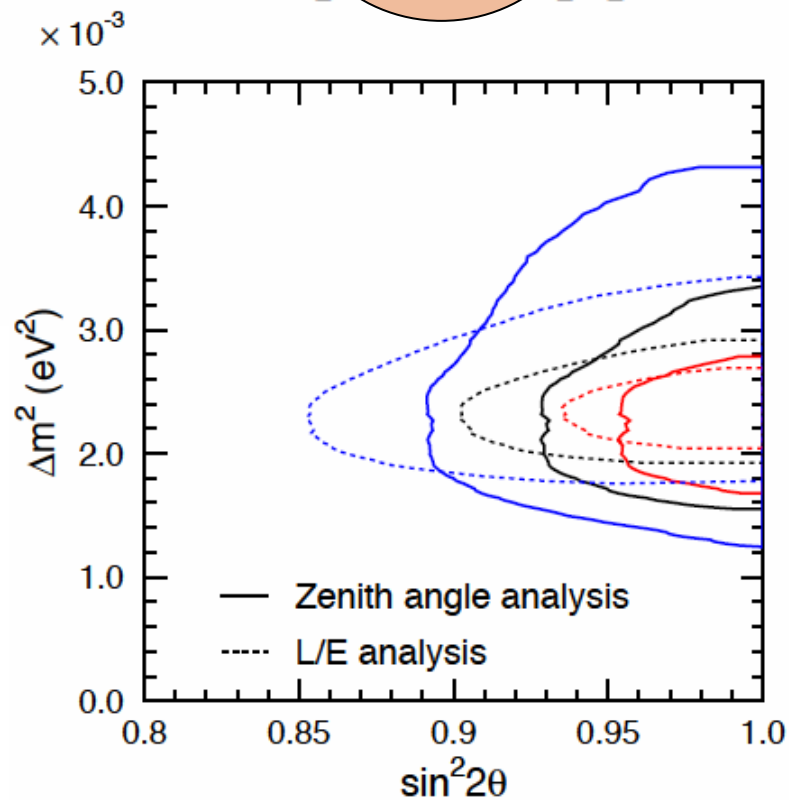


# $\nu$ oscillation has been seen!

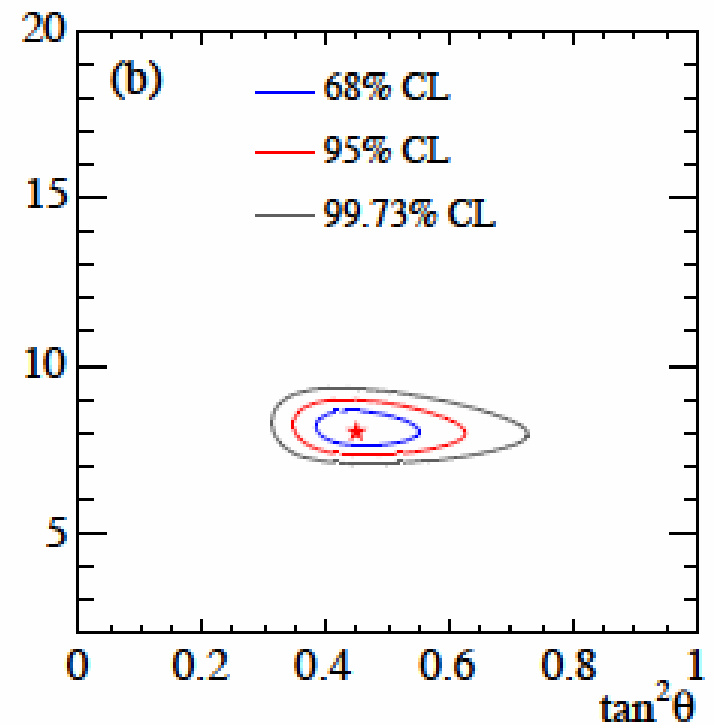


# MNS matrix and $\nu$ mass pattern

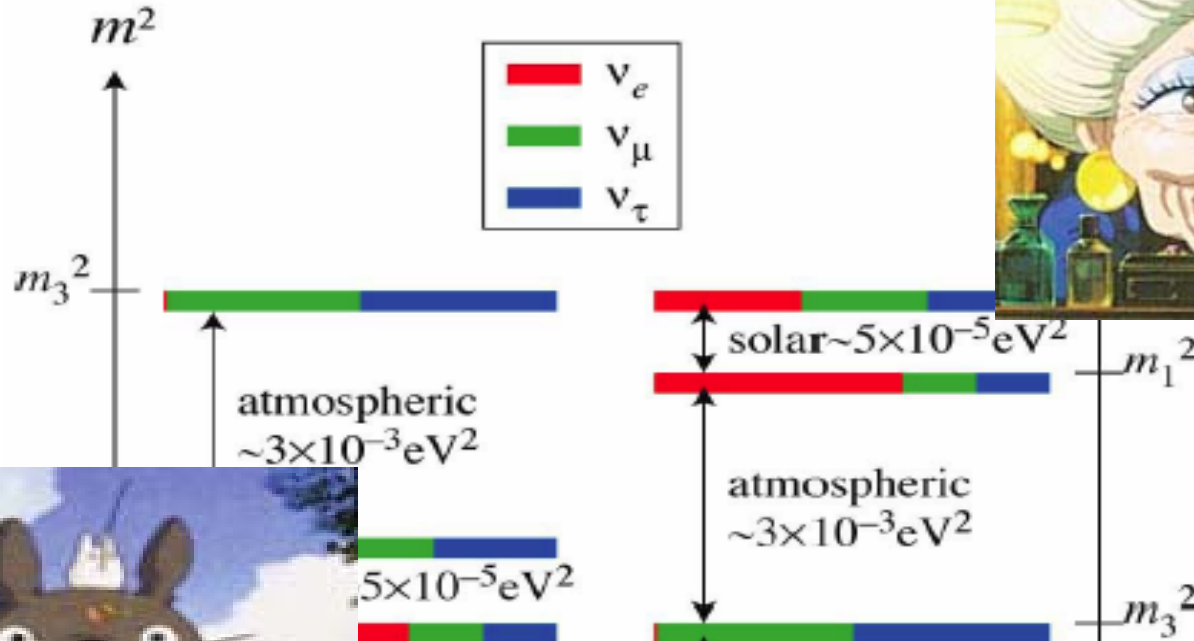
$$U \equiv U_{\text{MNS}} \cdot \Gamma = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \times \text{diag}(1, e^{i\beta}, e^{i\gamma})$$



$\Delta m^2 \text{ (} 10^{-5} \text{ eV}^2\text{)}$



# $\nu$ mass hierarchy & absolute mass scale



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# Pressing questions

- Origin of  $\nu$  masses and mixing
- Large lepton mixing vs. small quark mixing
- Quark lepton symmetry/relationship incl. flavor symmetry
- How to determine remaining parameters?...

Need for some strategic thoughts?

1. How to detect nonzero  $\theta_{13}$
2. How to measure CP violation phase  $\delta$
3. A coupled problem; CPV-mass hierarchy



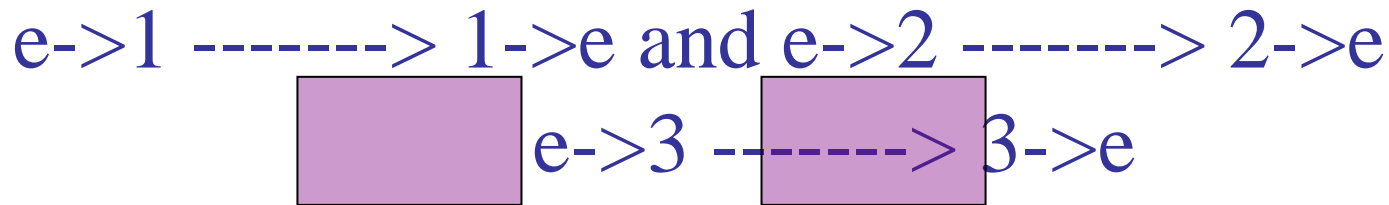
$\theta_{13}$  first

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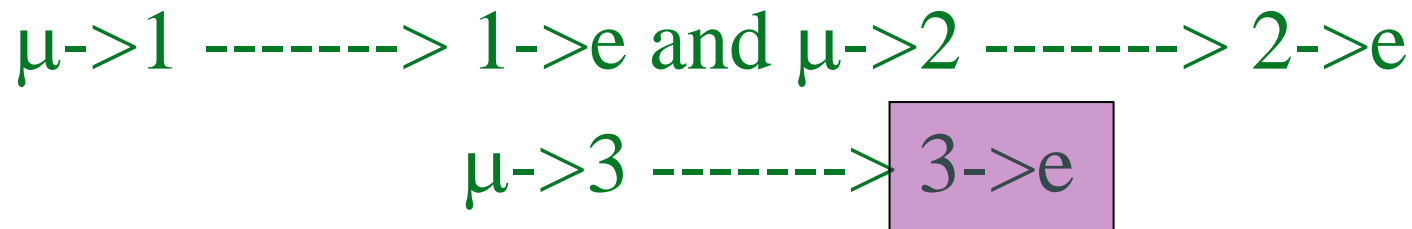
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To measure  $\theta_{13}$  one needs  $\nu_e$

- $P(\nu_e \rightarrow \nu_e)$  is the interference between



- $P(\nu_\mu \rightarrow \nu_e)$  is the interference between





# Reactor neutrino experiments

Reactor  $\bar{\nu}_e$

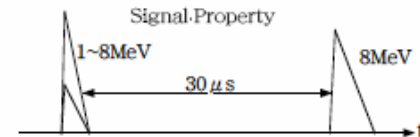
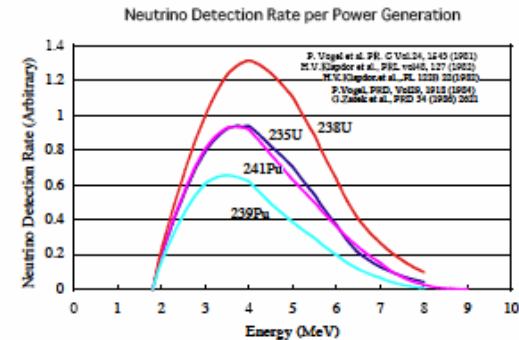
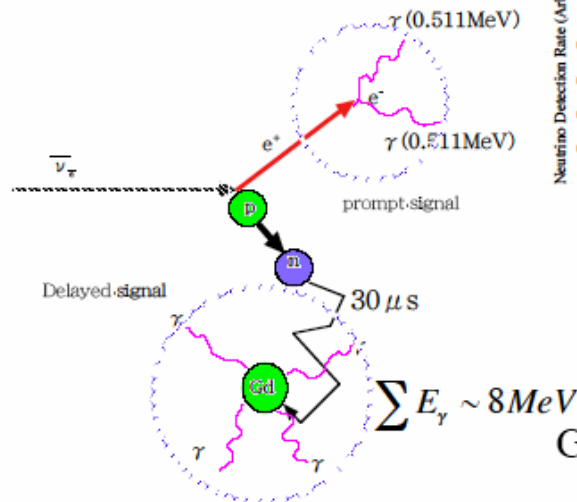
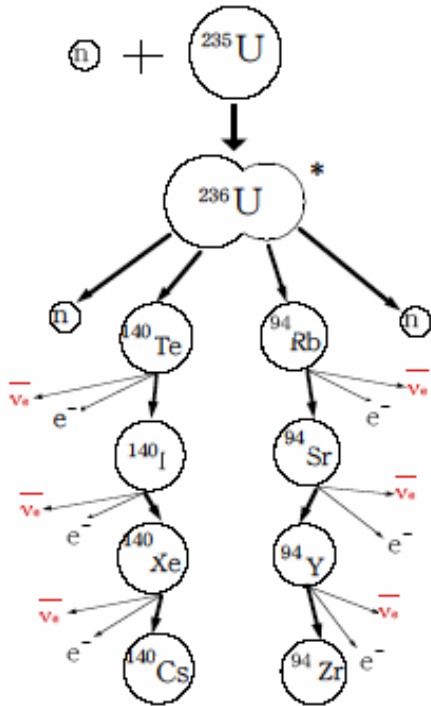
$\sim 6\nu/\text{fission}$  &  $\sim 200\text{MeV}/\text{fission}$



$\sim 6 \times 10^{20} \bar{\nu}_e / \text{s} / \text{reactor}$  (1GWe)

$\bar{\nu}_e$  Detection

(Most of the proposed projects uses Gd-LS)



Gd  $\Rightarrow$  largest  $n$  absorption  $\sigma$  & emits high energy  $\gamma$ s.

Delayed Coincidence

$\Rightarrow$  Background is severely suppressed

# Reactor measurement of $\theta_{13}$

$$1 - P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = \sin^2 2\theta_{13} \sin^2 \left( \frac{\Delta m_{31}^2 L}{4E} \right) + O(\epsilon s_{13}^2) + O(\epsilon^2)$$

$$\epsilon \equiv \Delta m_{21}^2 / \Delta m_{31}^2 \simeq 0.03$$

- Independent of  $\delta$ , matter effect,  $\theta_{23}$ ,  $\theta_{12}$ , solar  $\Delta m^2$

$\Rightarrow$  Pure measurement of  $\theta_{13}$

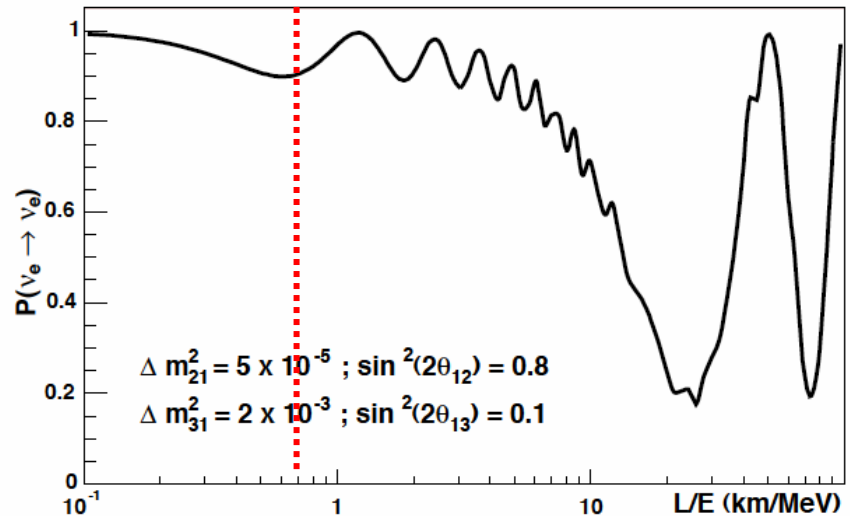
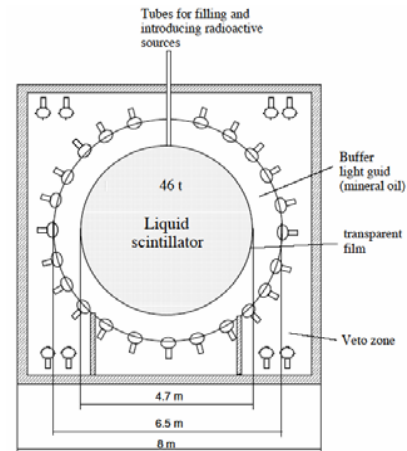


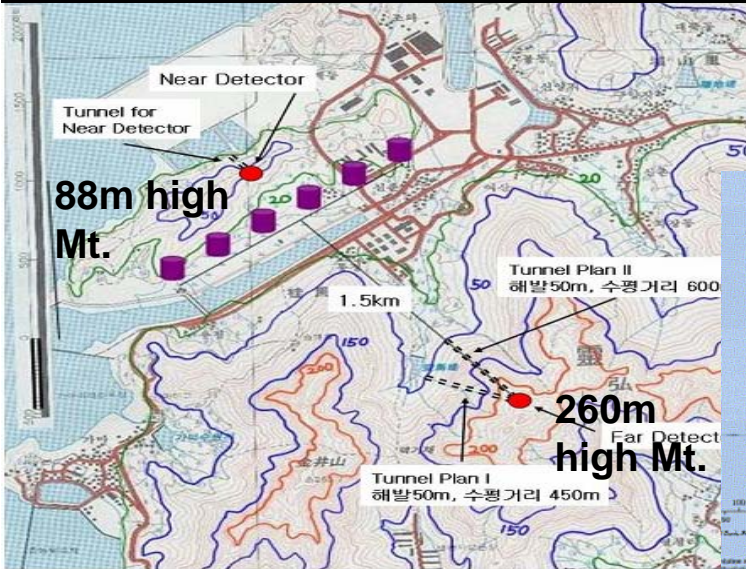
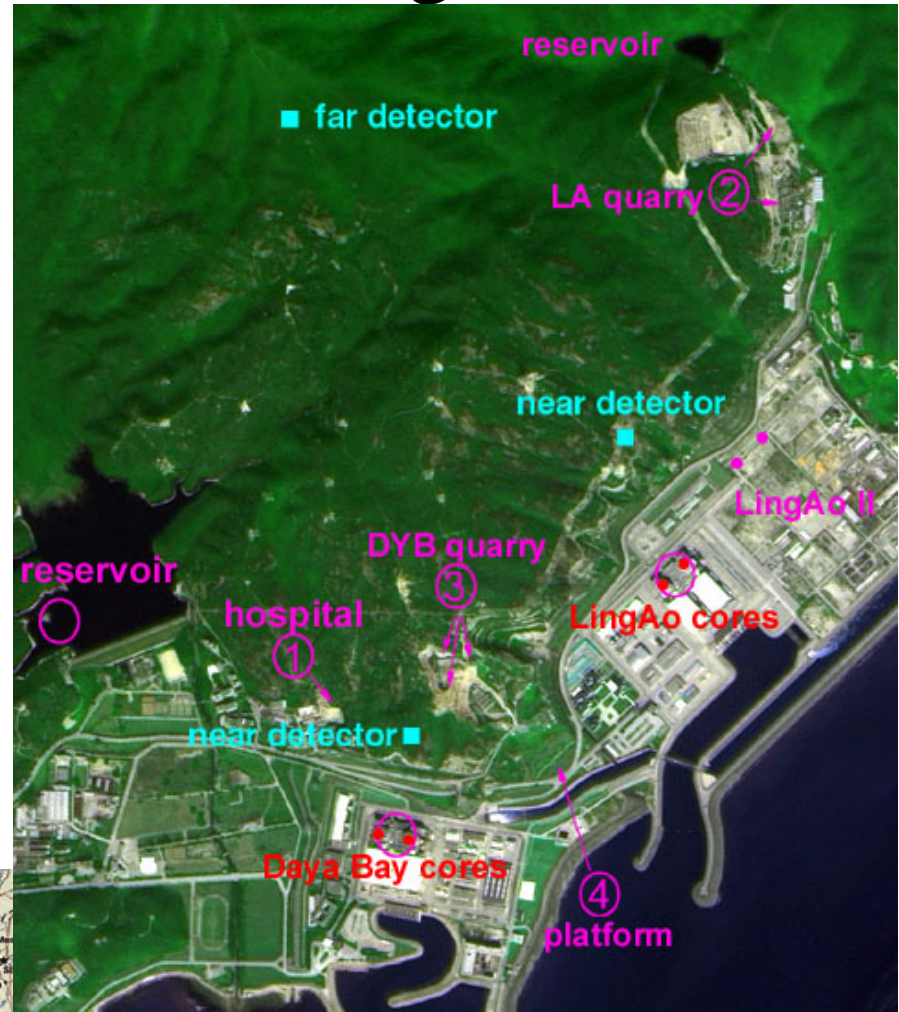
Figure 3: Probability of  $\nu_e$  disappearance versus  $L/E$  for  $\theta_{13}$  at its current upper limit

# Varying proposals over the globe

## The Chooz site



2x12.5 tons, D1=100-200m, D2=1050m. Sensitivity: 3 years  $\rightarrow \sin^2(2\theta_{13}) < \sim 0.03$



# LBL measurement of $\theta_{13}$

$$P(\nu_\mu \rightarrow \nu_e) = |\sqrt{P_{atm}} + e^{i(\delta \pm \frac{\Delta_{31}}{2})} \sqrt{P_{solar}}|^2$$

$$P_{atm} = \left( s_{13} s_{23} \Delta_{31} \frac{\sin\left(\frac{\Delta_{31} \mp aL}{2}\right)}{\left(\frac{\Delta_{31} \mp aL}{2}\right)} \right)^2$$

$$P_{solar} = \left( c_{12} s_{12} c_{23} \Delta_{21} \frac{\sin\left(\frac{aL}{2}\right)}{\left(\frac{aL}{2}\right)} \right)^2$$

$$\Delta_{31} \equiv \frac{|\Delta m_{31}^2| L}{2E}, \quad a = \sqrt{2} G_F N_e(x),$$

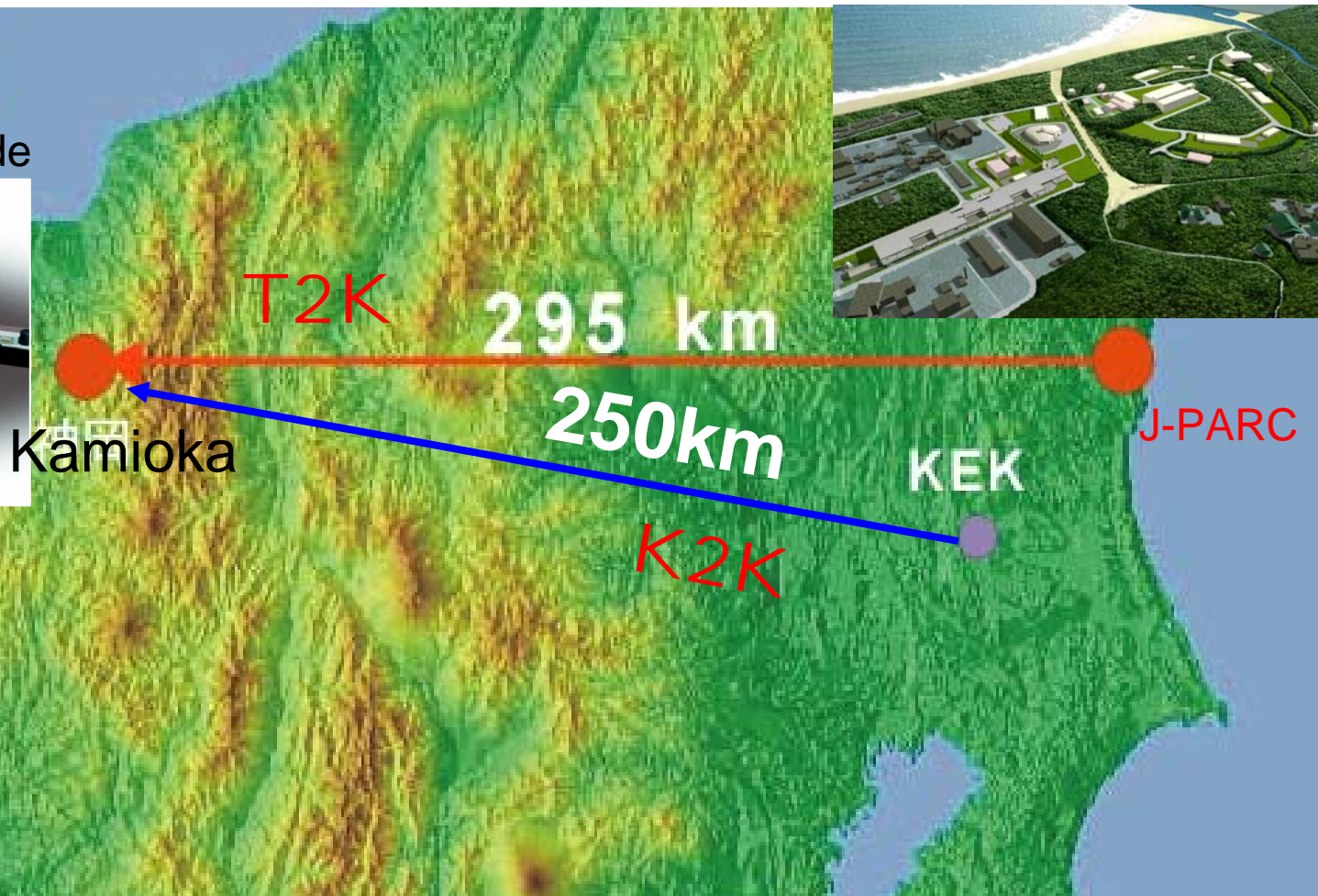
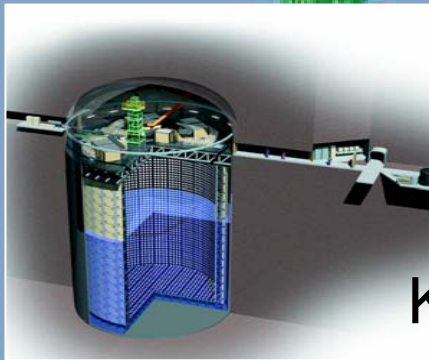
$$\pm = \text{sign of } \Delta m_{31}^2$$

Some slides come from:

Takaaki Kajita  
@NOW2006  
(Otoranto, Italy)  
for T2K

# T2K Phase-I

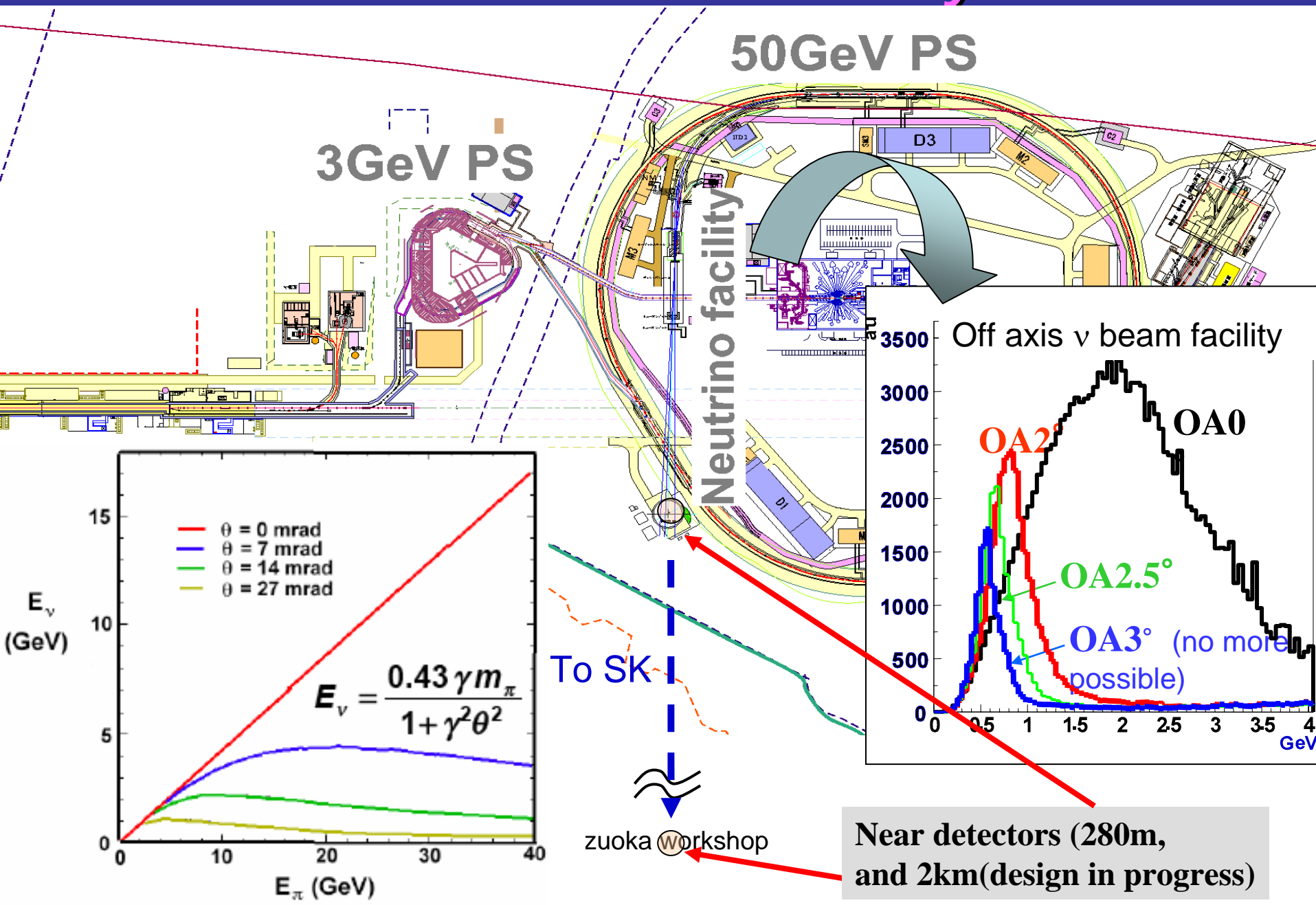
Super-Kamiokande



Collaboration (at present): Canada, France, Germany, Italy, Japan, Korea, Poland, Russia, Spain, Switzerland, UK, USA

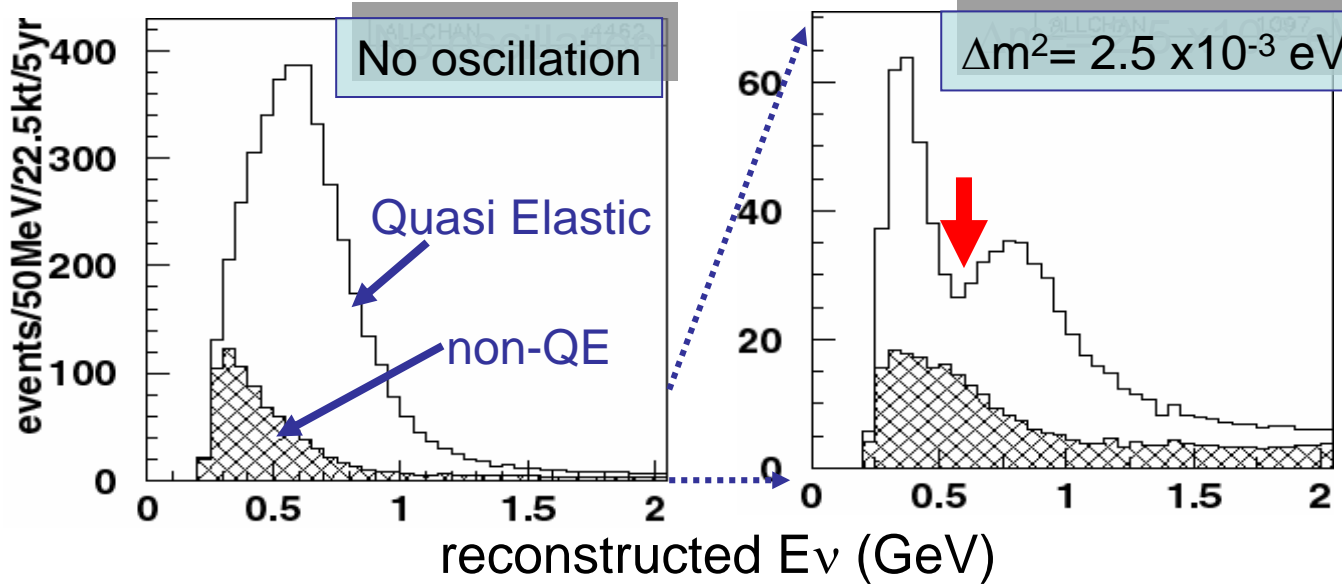
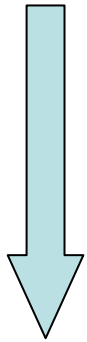
For details: I. Kato' talk

# T2K neutrino facility



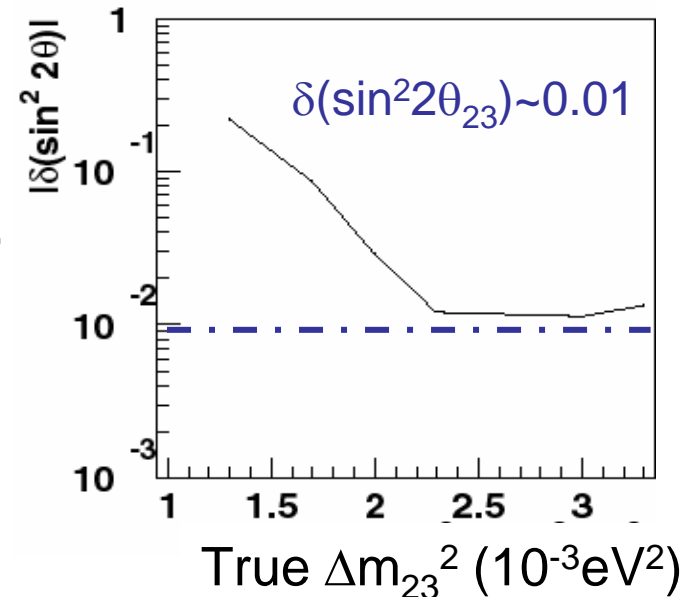
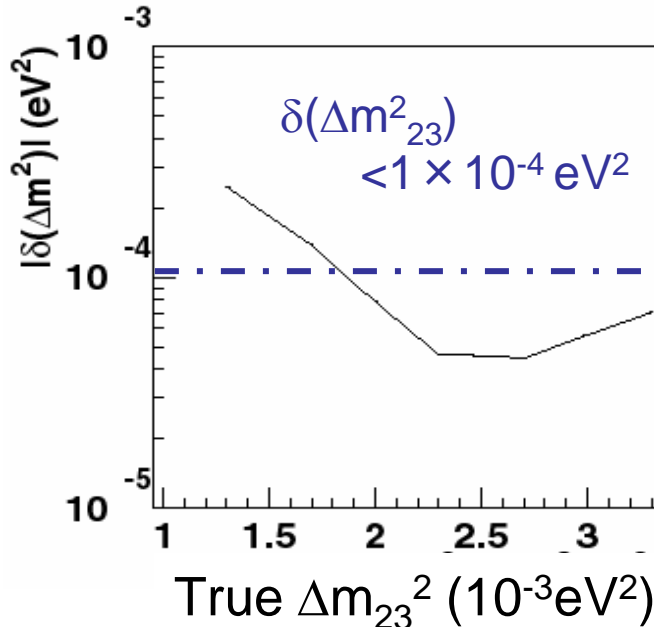
# Measurement of $\Delta m^2$ and $\sin^2 2\theta_{23}$

Reconstructed  $\nu$  energy distribution for single ring  $\mu$ -events.



Sensitivity with 0.75MW·5years and

- normalization ( 5%)
- non-qe/qe ratio ( 20%)
- E scale ( 1%)
- Spectrum shape(5%)
- Spectrum width ( 5%)



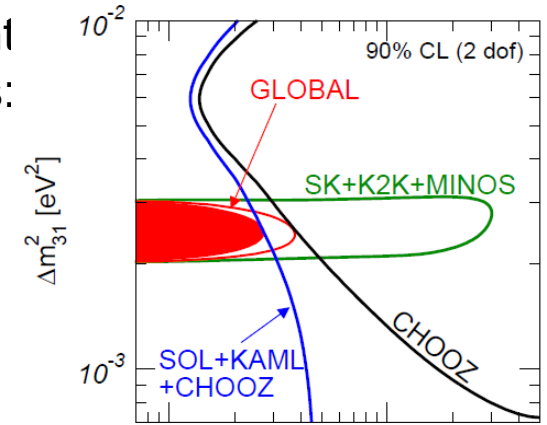
# $\theta_{13}$

$$P(\nu_{\mu} \rightarrow \nu_e)$$

$$= \sin^2 \theta_{23} \cdot \sin^2 2\theta_{13} \cdot \sin^2 \left( \frac{1.27 \Delta m_{23}^2 L}{E} \right)$$

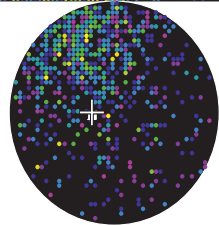
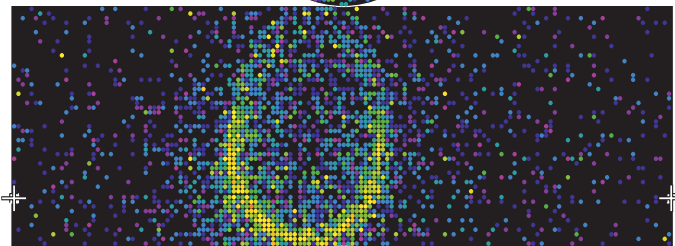
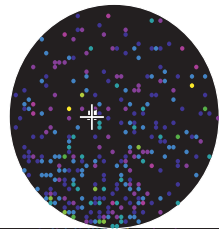
( $\Delta m_{12}^2=0$  assumed, matter effect not included)

Present status:

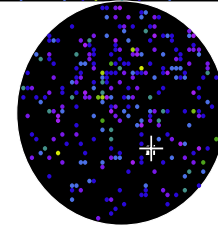
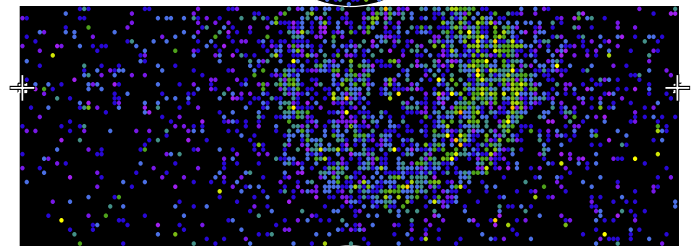
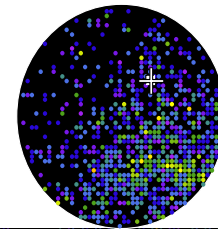


$\sin^2 \theta_{13}$  T.Schwetz  
hep-ph/0606060

Signal for non-zero  $\theta_{13}$  ( $\nu_{\mu} \rightarrow \nu_e$ )



BG (NC  $\pi^0$  ....)

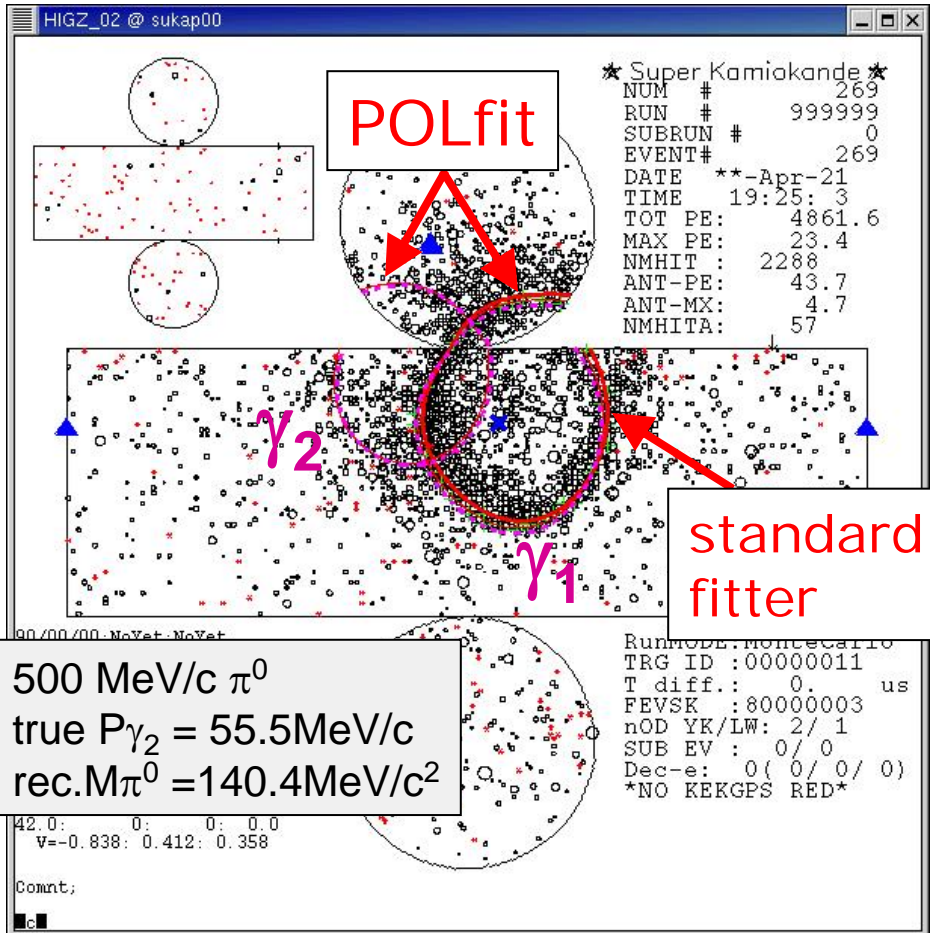


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# POL(Pattern of Light)fit – $\pi^0$ fitter –



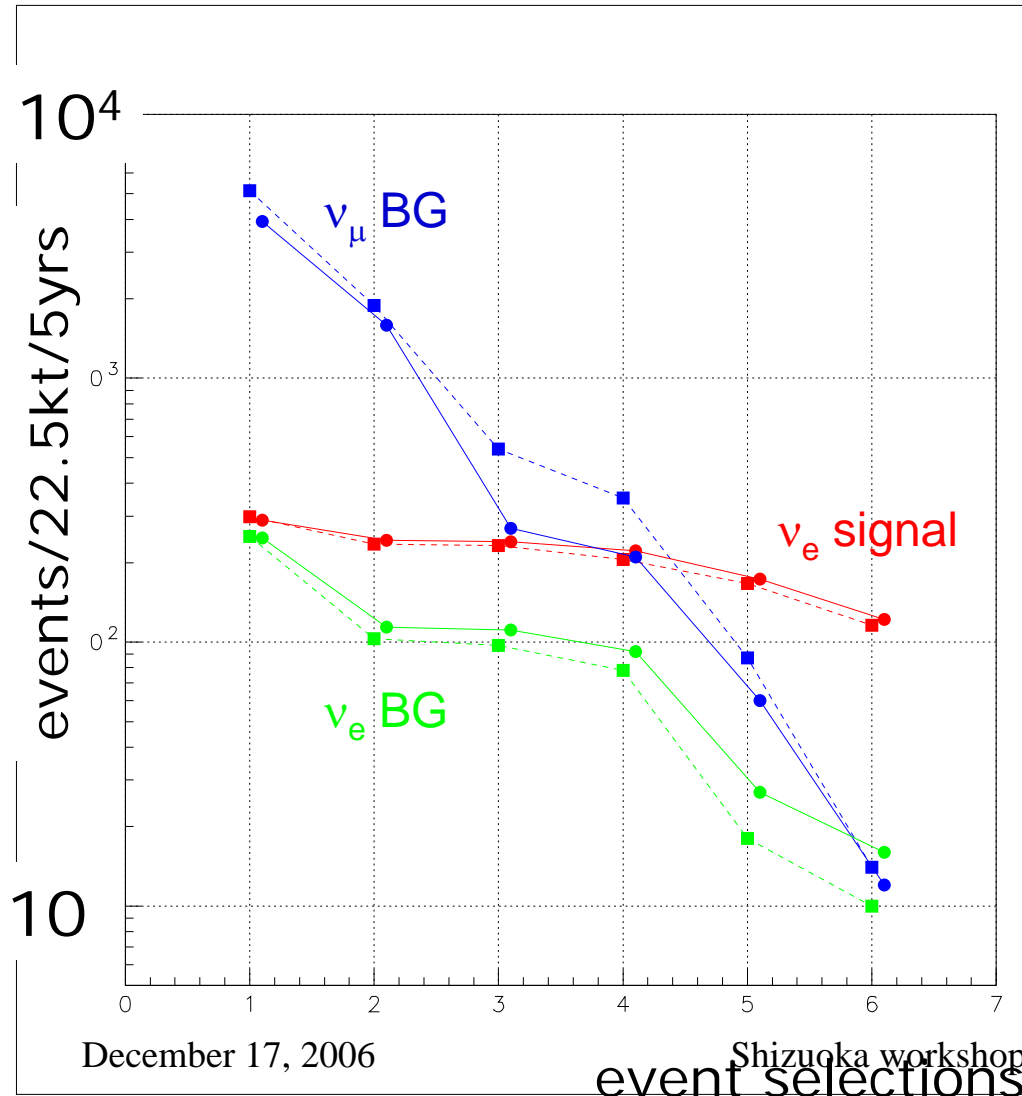
- Target: FCFV 1R-like events
- $\Delta L \equiv \text{Likelihood}(2\gamma \text{ assump.}) - \text{Likelihood}(\text{electron assump.})$
- Try to reconstruct two  $\gamma$  rings
- Input: vertex, visible energy, and the 1<sup>st</sup>  $\gamma$  direction by the standard fitter
- Compare observed & expected (direct+scatter) charge
- Vary the 2<sup>nd</sup>  $\gamma$  direction and the energy fraction until the best match found

→  $M\pi^0$  etc.

# Events vs. Selections

$$\Delta m^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{13} = 0.1$$



Event selections:

1. FCFV,  $E_{\text{vis}} > 100$
2. 1 ring
3. e-like
4. no decay-e
5.  $0.35 < E_{\nu}^{\text{rec}} < 0.85$
6.  $\pi^0$  cuts:
  - $\cos < 0.90$
  - $\Delta L < 80, M_{\pi^0} < 100$

# Events vs. selections

$$\Delta m^2 = 2.5 \times 10^{-3} \text{eV}^2, \sin^2 2\theta_{13} = 0.1$$

(events / 22.5kt / 5yrs)

	$\nu_\mu$ CC BG	$\nu_\mu$ NC BG	beam $\nu_e$ BG	$\nu_e$ (CC) Signal
FCFV, $E_{\text{vis}} > 100$	<b>2849</b>	<b>1082</b>	<b>248</b>	<b>290</b>
1R	<b>1313(46%)</b>	<b>277(26%)</b>	<b>114(46%)</b>	<b>243(84%)</b>
e-like	<b>51(1.8%)</b>	<b>219(20%)</b>	<b>111(45%)</b>	<b>240(83%)</b>
no decay-e	<b>15(0.5%)</b>	<b>195(18%)</b>	<b>92(37%)</b>	<b>222(77%)</b>
$0.35 < E_v^{\text{rec}} < 0.85$	<b>2.2(0.1%)</b>	<b>58(5%)</b>	<b>27(11%)</b>	<b>173(60%)</b>
$\Delta L < 80, M < 100, \cos < 0.9$	<b><math>12 \pm 0.8(0.3\%)</math></b> (stat.)		<b><math>16 \pm 0.4(6\%)</math></b> (stat.)	<b><math>122 \pm 3(42\%)</math></b> (stat.)



(old  $\pi^0$  fitter:

12

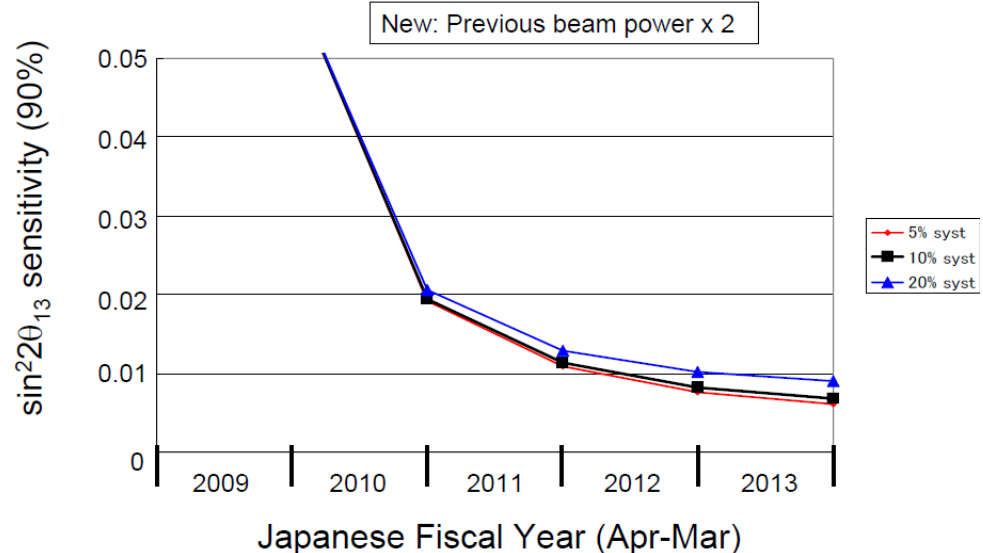
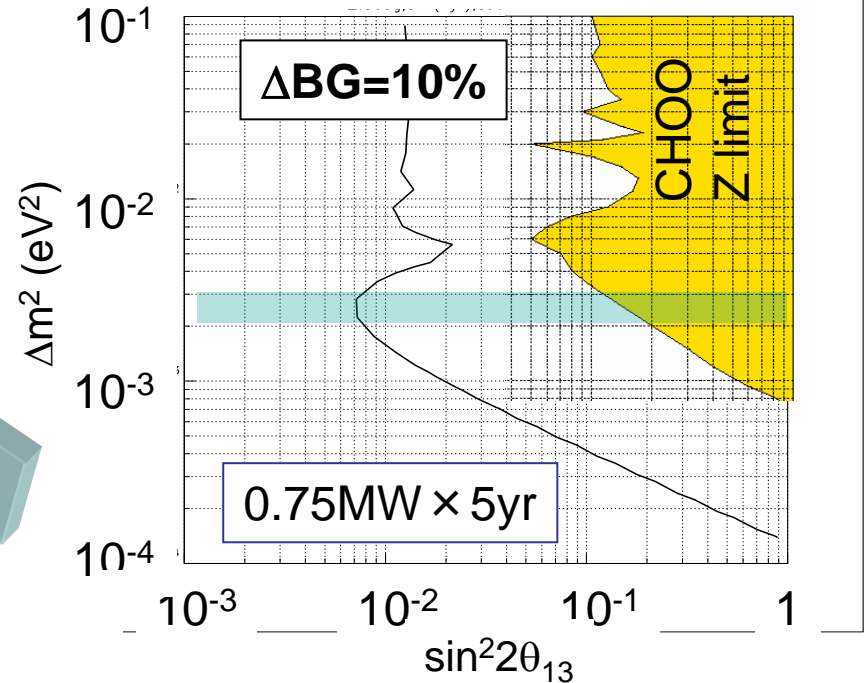
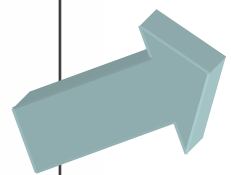
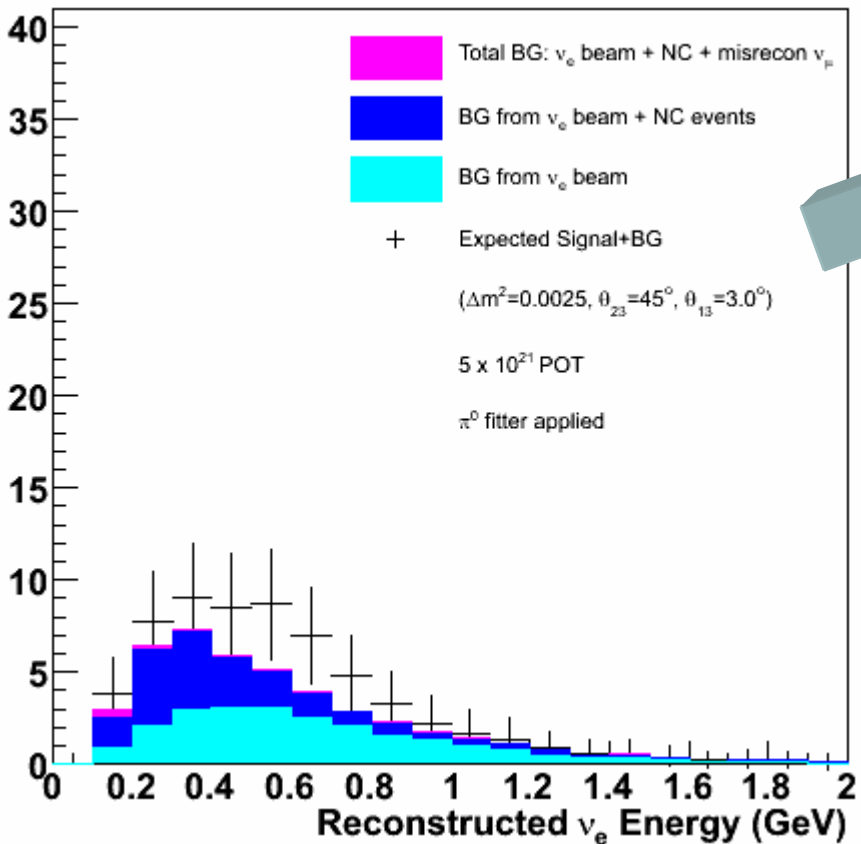
15

109)

# $\nu_e$ appearance and $\theta_{13}$ sensitivity

$\nu_e$  appearance signal and BG at SK.

$$\sin^2 2\theta_{13} = 0.01$$

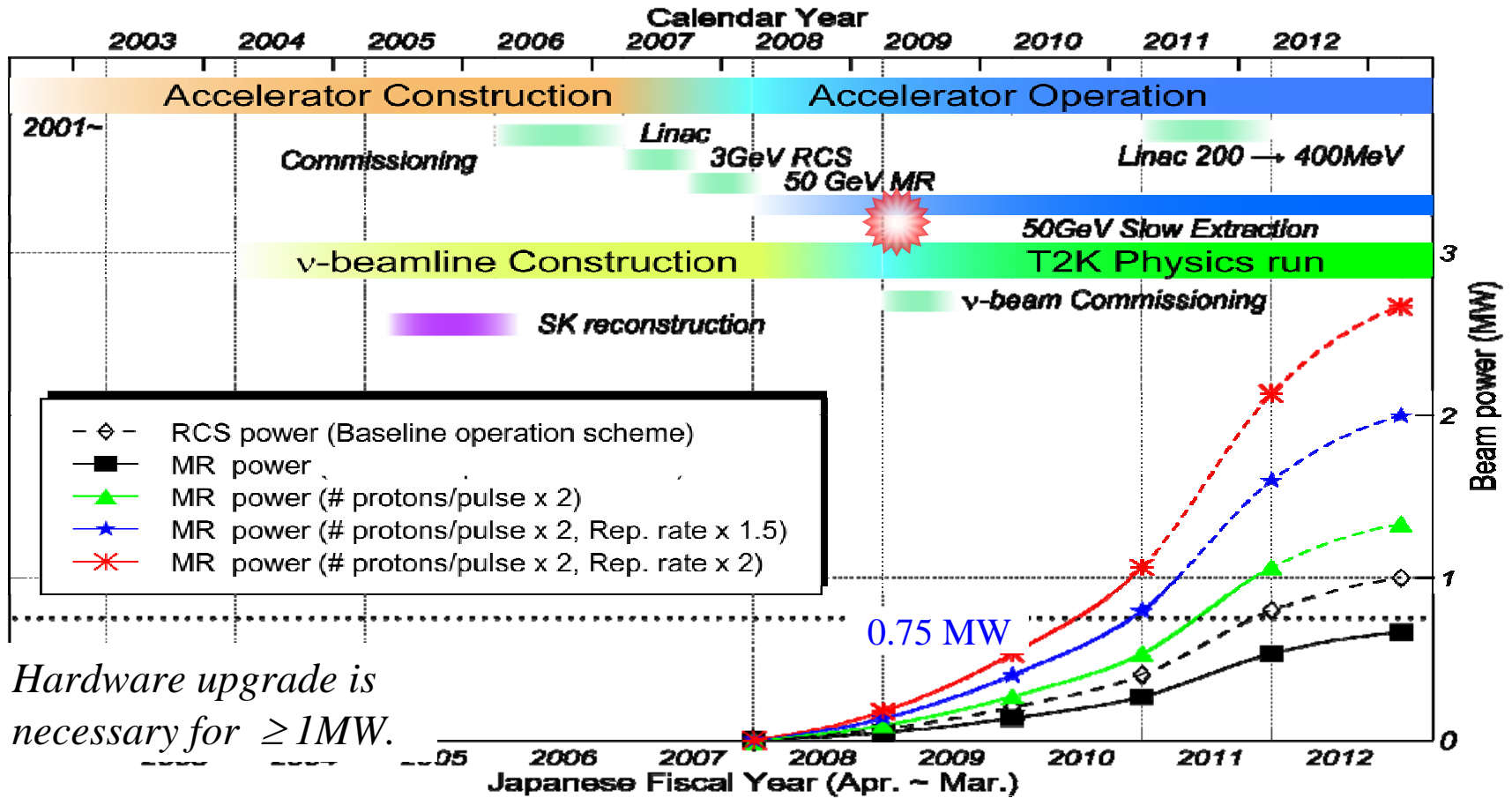


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# J-PARC schedule & Beam Power estimation

Nakadaira neutrino2006



Target date for new RF system installation.

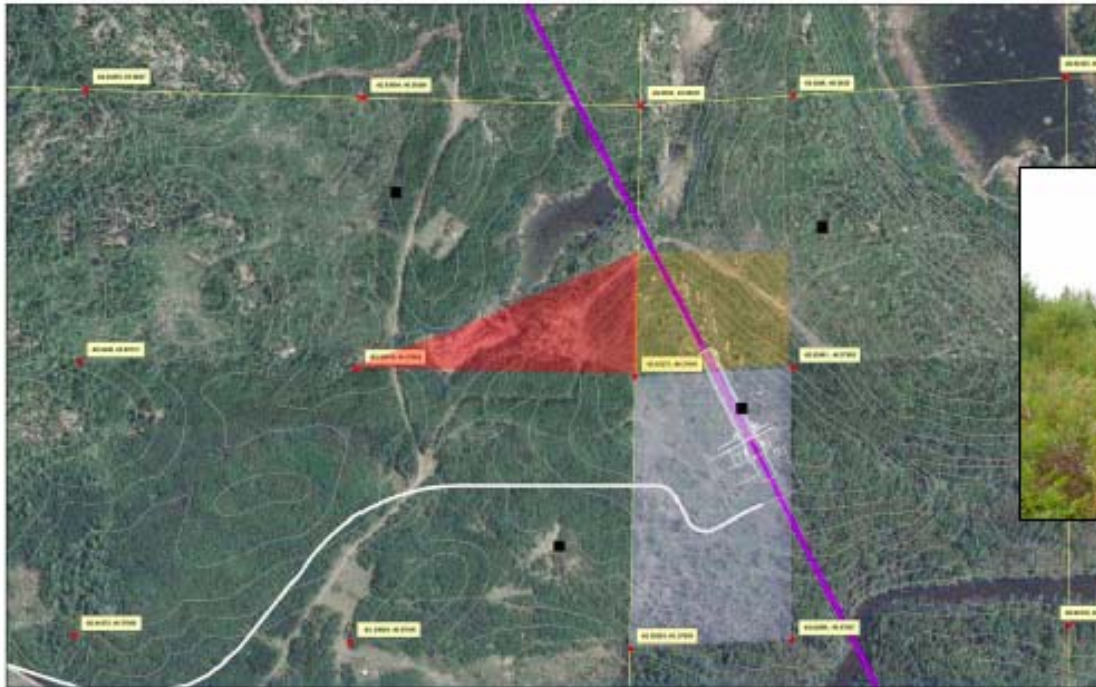
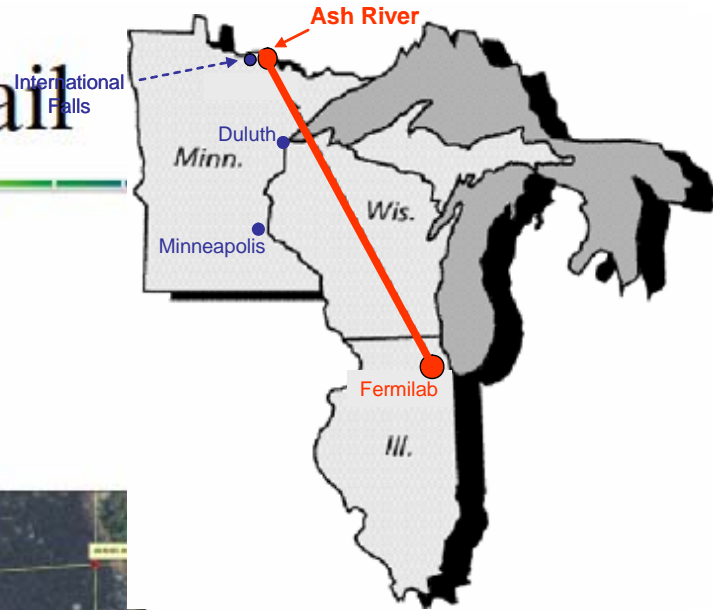
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# Site – Ash River Trail

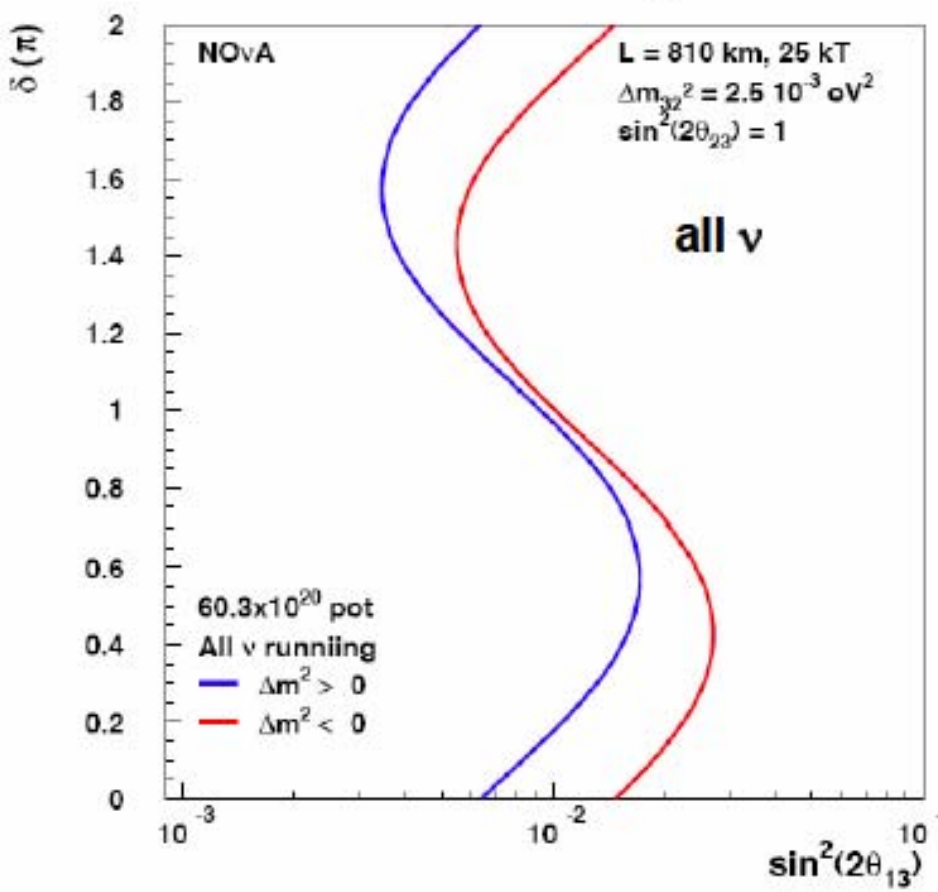
- 503 miles (810 km) from Fermilab
- 3.6 Mile Access Road
- Electrical Upgrade



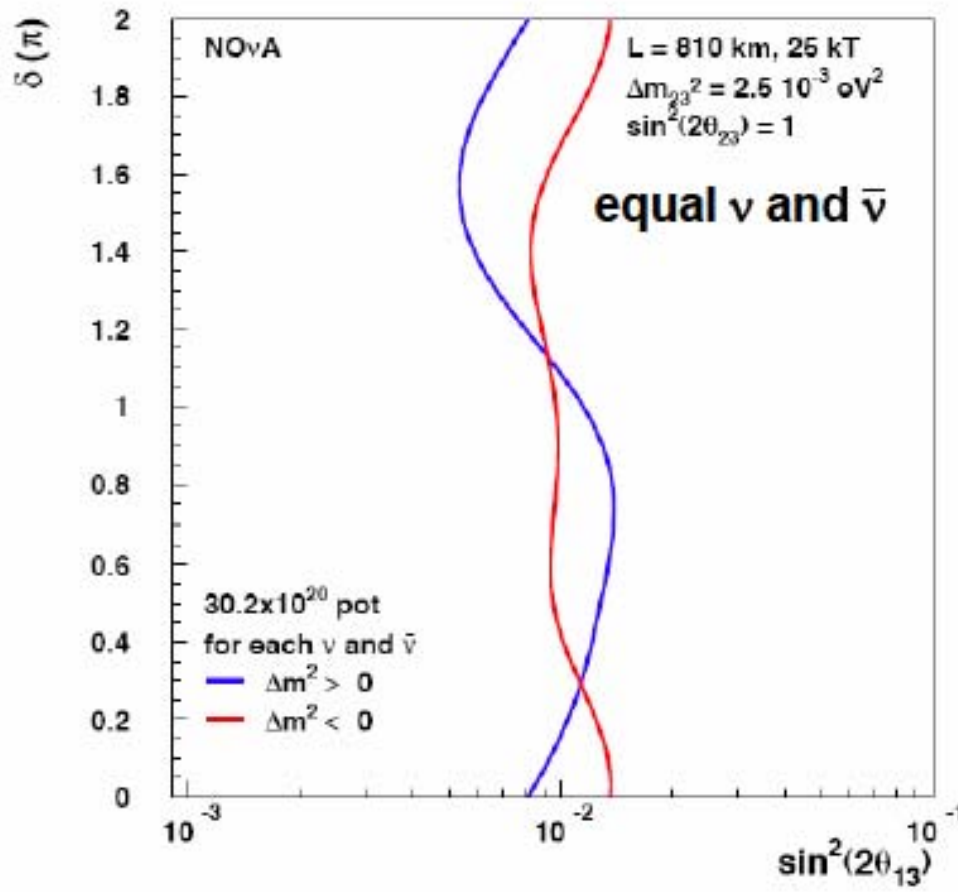


# 3 $\sigma$ Sensitivity to $\theta_{13} \neq 0$

3  $\sigma$  Sensitivity to  $\sin^2(2\theta_{13}) \neq 0$



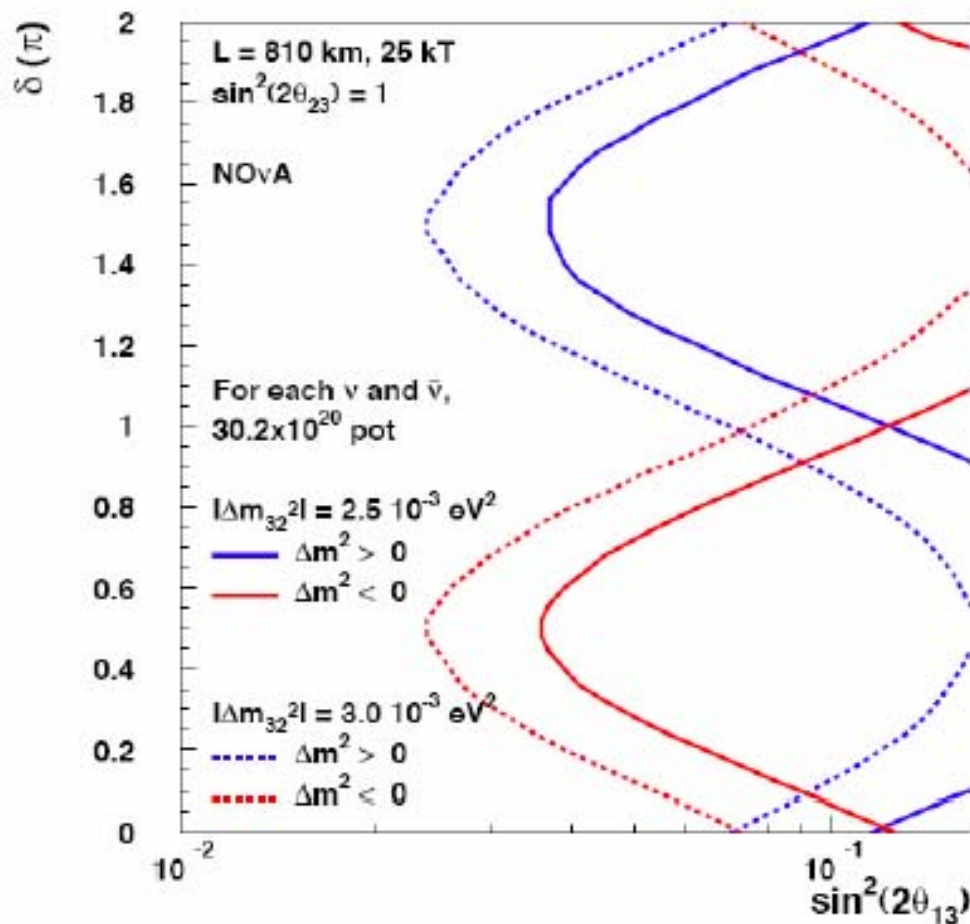
3  $\sigma$  Sensitivity to  $\sin^2(2\theta_{13}) \neq 0$





# 95% CL Resolution of the Mass Ordering

95% CL Resolution of the Mass Hierarchy



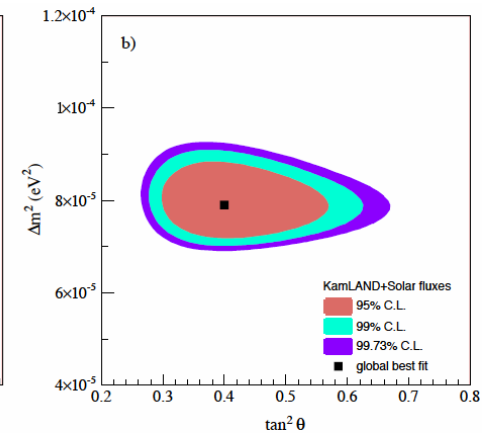
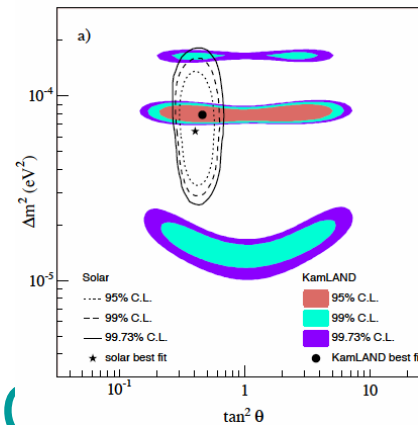


# Need for beyond the next generation (NG) experiments

- NG exp. will not determine (unless very lucky) the mass hierarchy
- NG exp. does not have sensitivity to CP violation
- NG exp. may not be able to see nonzero  $\theta_{13}$  (what happens then?)
- Question: how accurately should we need to know  $\Delta m^2$  and  $\theta$ 's?

# Quark-lepton complementarity ?

$$\theta_C + \theta_{\text{solar}} = 45.1^\circ \pm 2.4^\circ (1\sigma)$$



$$36.8^\circ < \theta_{\text{atm}} < 53.2^\circ (90\% \text{ C.L.})$$

$$2.3^\circ < \theta_{23}^q < 2.5^\circ (90\% \text{ C.L.})$$

$$\theta_{23}^q + \theta_{\text{atm}} = 47.4^\circ \pm 8.3^\circ (90\% \text{ C.L.})$$

# Foreseeing the future

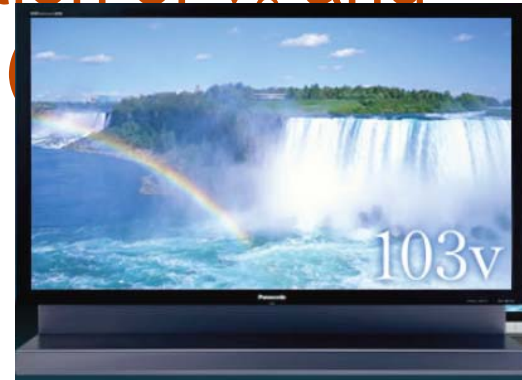


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# Things changes at $\sin^2 2\theta_{13} \sim 0.01$

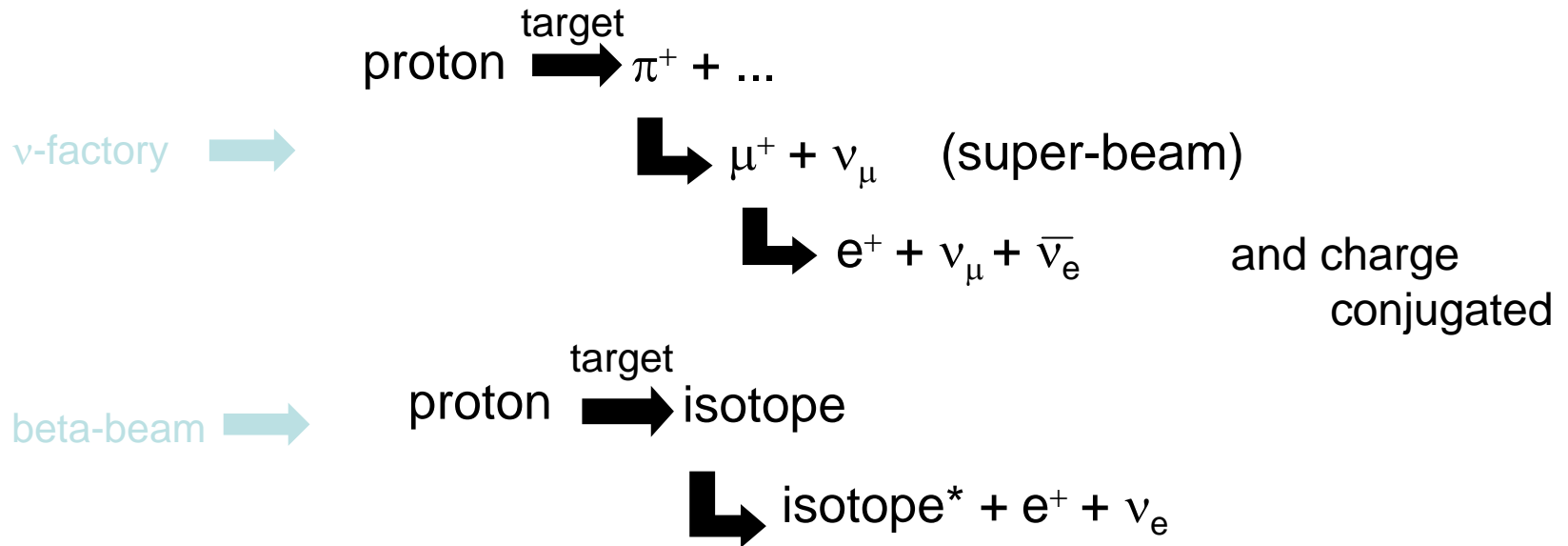
- Conventional super  $\nu_\mu$  beam + Mton water detector work
- Known beam technology
- Background highly nontrivial
- $\nu_e$  beam contamination not negligible but tolerable
- beta beam / neutrino factory required
- Requires long-term R&D efforts
- Low background
- pure  $\nu_e$  beam ( $\beta$ ) / well understood combination of  $\nu_e$  and  $\nu_\mu$  beam



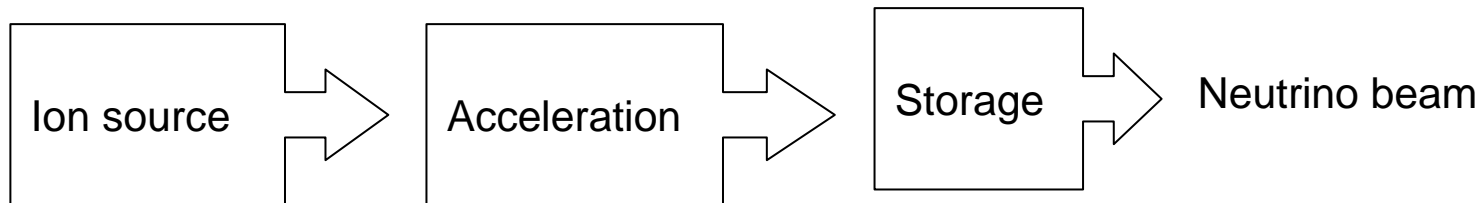
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# Various beam options



- v-factory uses beam of 4<sup>th</sup> generation.
- Beta-beam uses 3<sup>rd</sup> generation beam.
- Beta-beam is technically closer to existing/used accelerator technology.



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# Degeneracy; a notorious obstacle



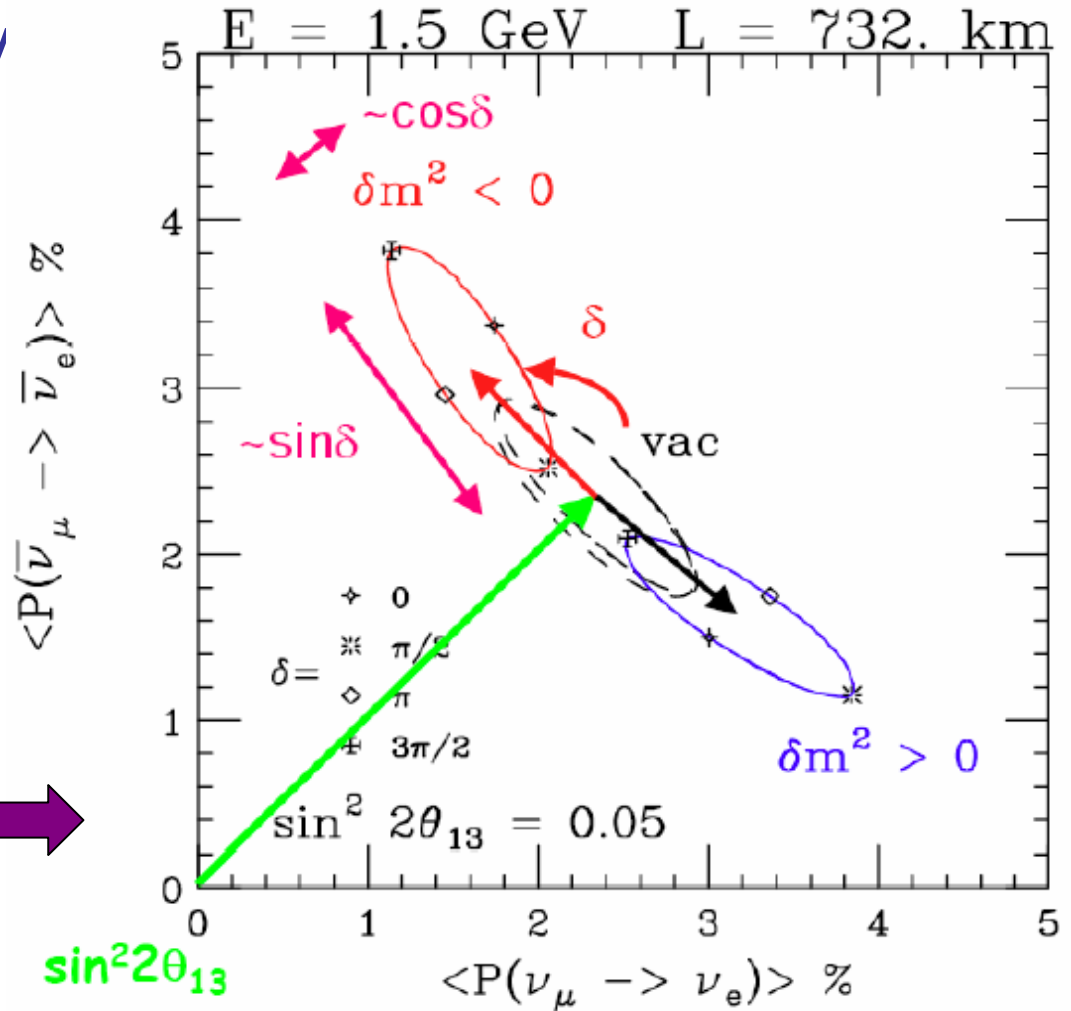
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# A machinery in my talk

Oscillation probability  
draw ellipse if  
plotted in bi-P  
plane

Role played by CP  
phase  $\delta$  and the  
matter clearly  
distinguished



# Cause of the degeneracy; easy to understand

- You can draw two ellipses from a point in P-Pbar space



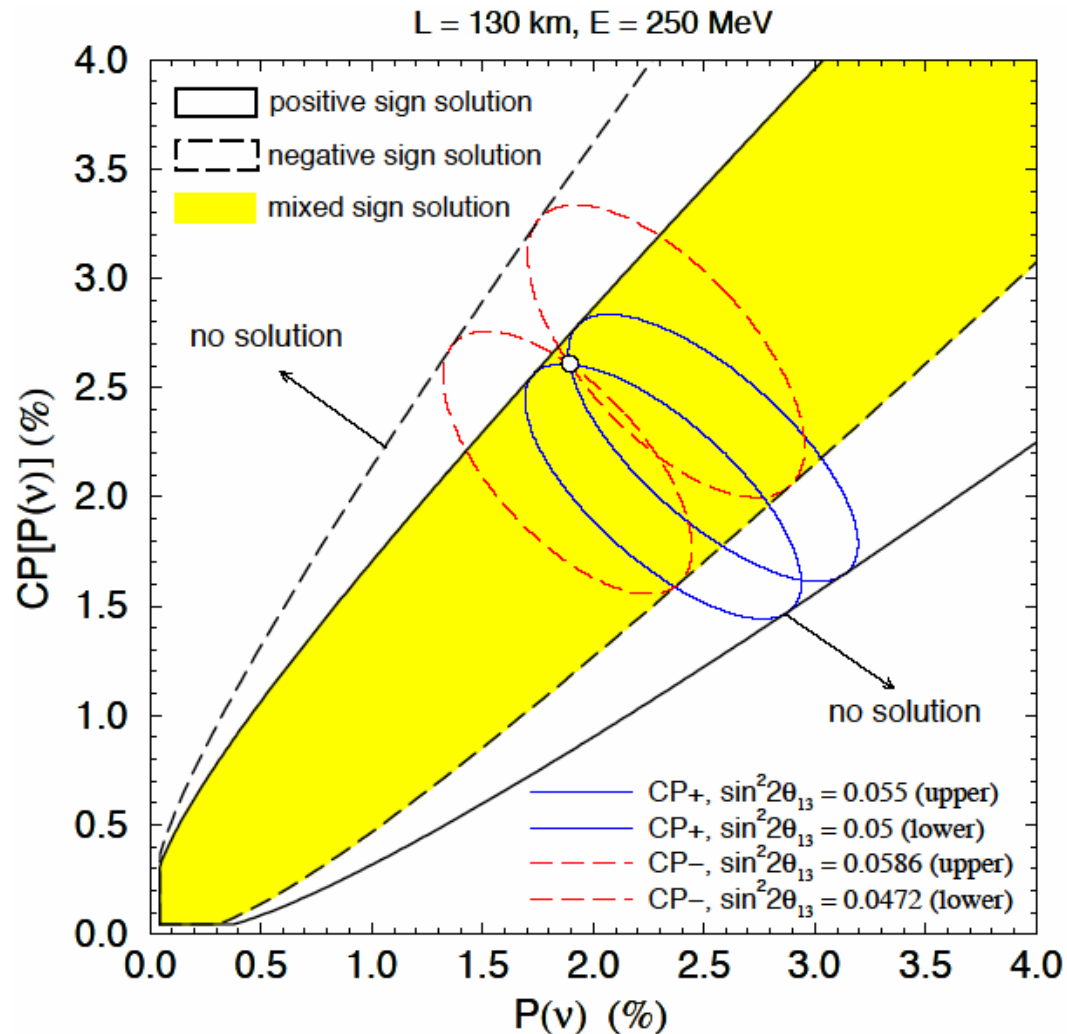
Intrinsic degeneracy

- Doubled by the unknown sign of  $\Delta m^2$



4-fold degeneracy

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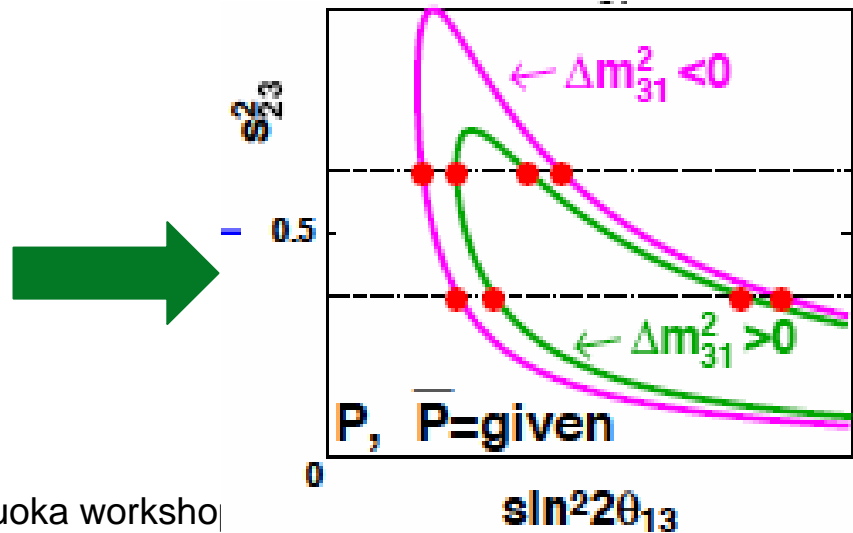
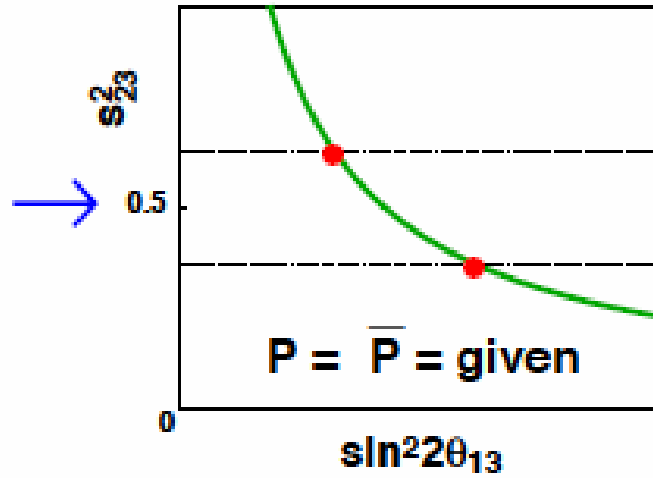
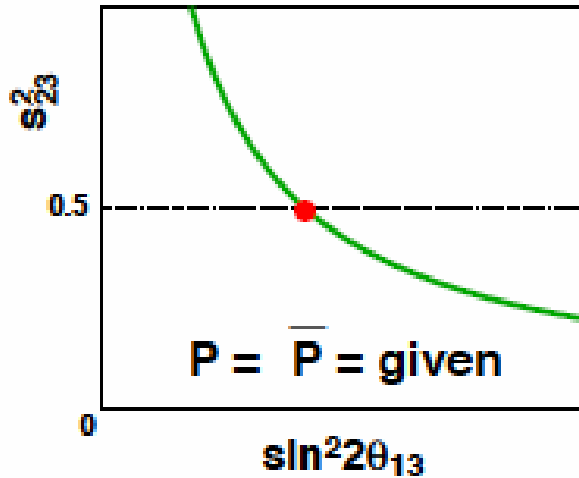
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# $\theta_{23}$ octant degeneracy

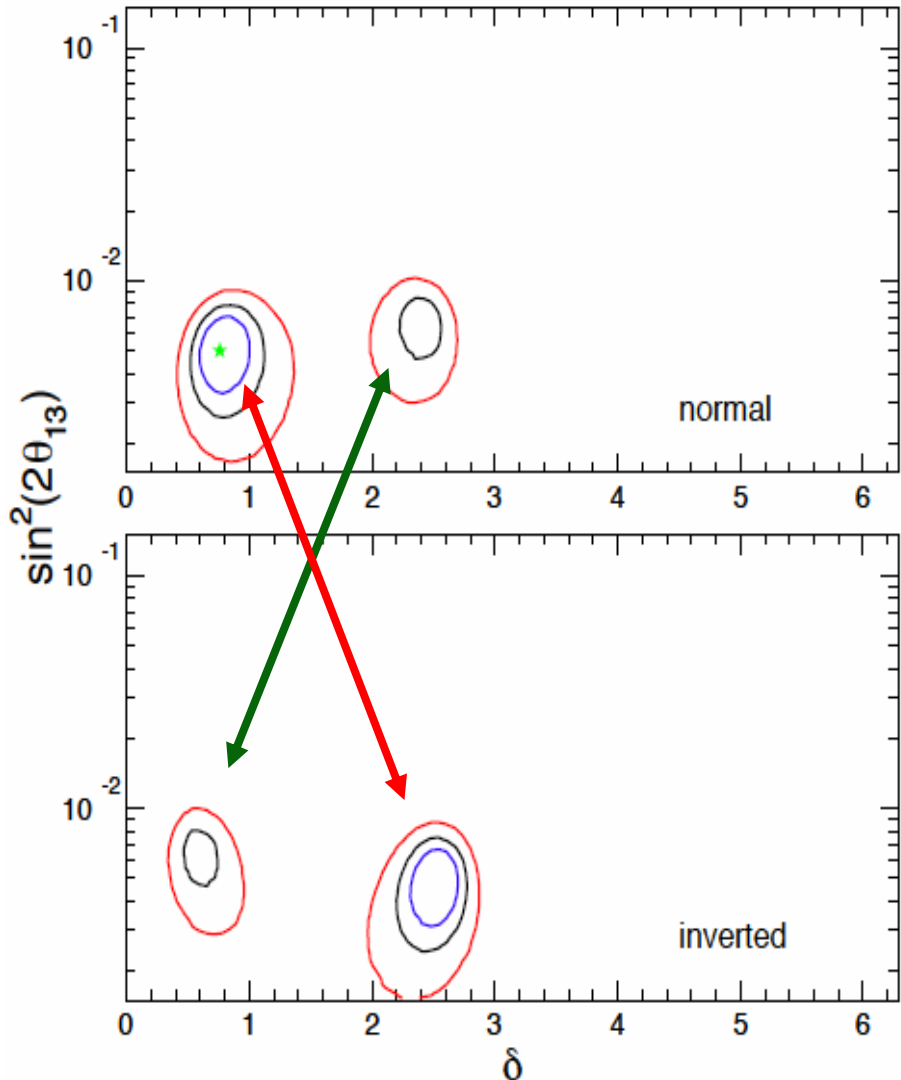
(a)  $\theta_{23} = \frac{\pi}{4}$ ,  $\Delta m_{21}^2 = 0$ ,  $A = 0$

(b)  $\theta_{23} \neq \frac{\pi}{4}$ ,  $\Delta m_{21}^2 = 0$ ,  $A = 0$



# Structure of intrinsic & sign- $\Delta m^2$ degeneracy in (matter) perturbative regime

(Kamioka 1Mt)  $\times$  (4MW,  $\nu$  2yr +  $\bar{\nu}$  6yr)

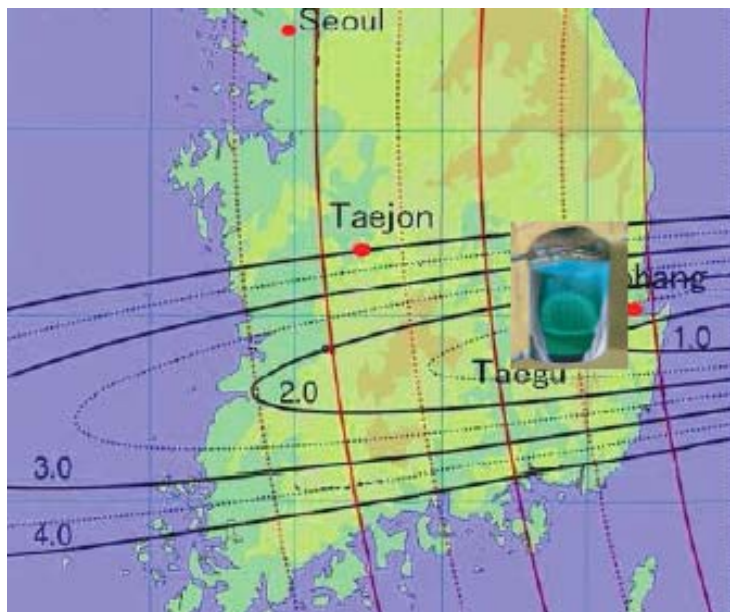


- Intrinsic degeneracy;  
 $\delta_2 = \pi - \delta_1$
- sign( $\Delta m^2$ )- $\delta$  degeneracy  
arises because  $P$  is  
approx. invariant under:

- $\Delta m^2 \longrightarrow -\Delta m^2$
- $\delta \xrightarrow{\hspace{1cm}} \pi - \delta$

# T2KK; Tokai-to-Kamioka-Korea identical two-detector complex

- An improvement over T2K II design with Hyper-K @ Kamioka with 1 megaton water



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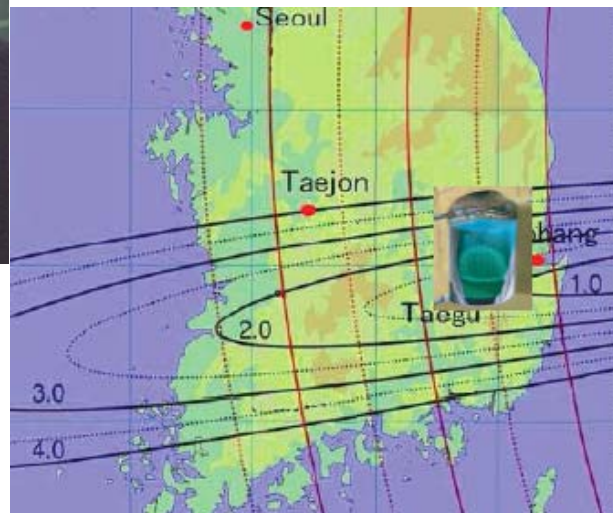
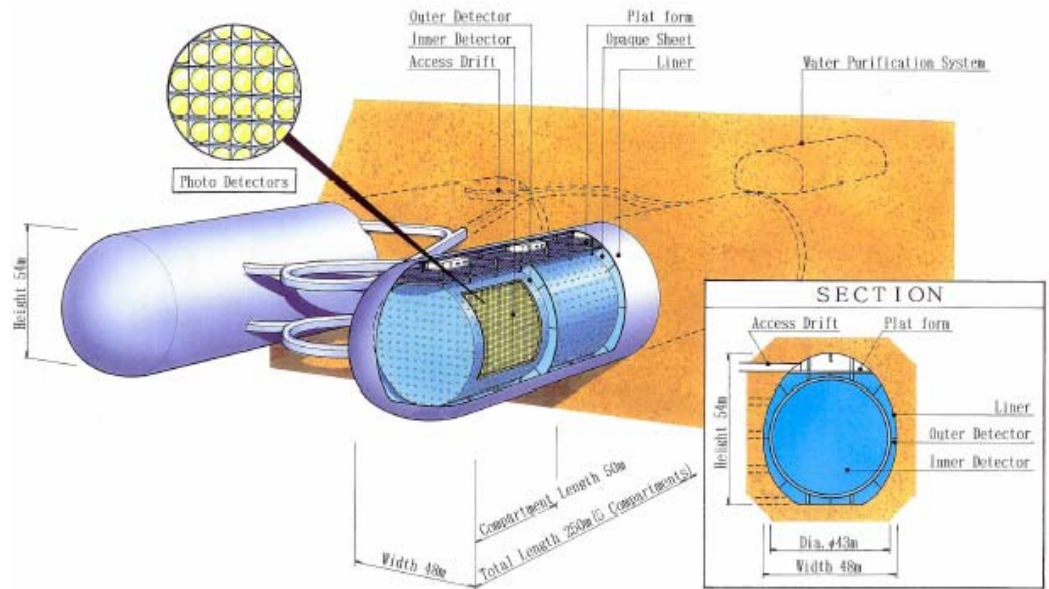


What's  
good in  
T2KK?  
(what about  
NOVA?)

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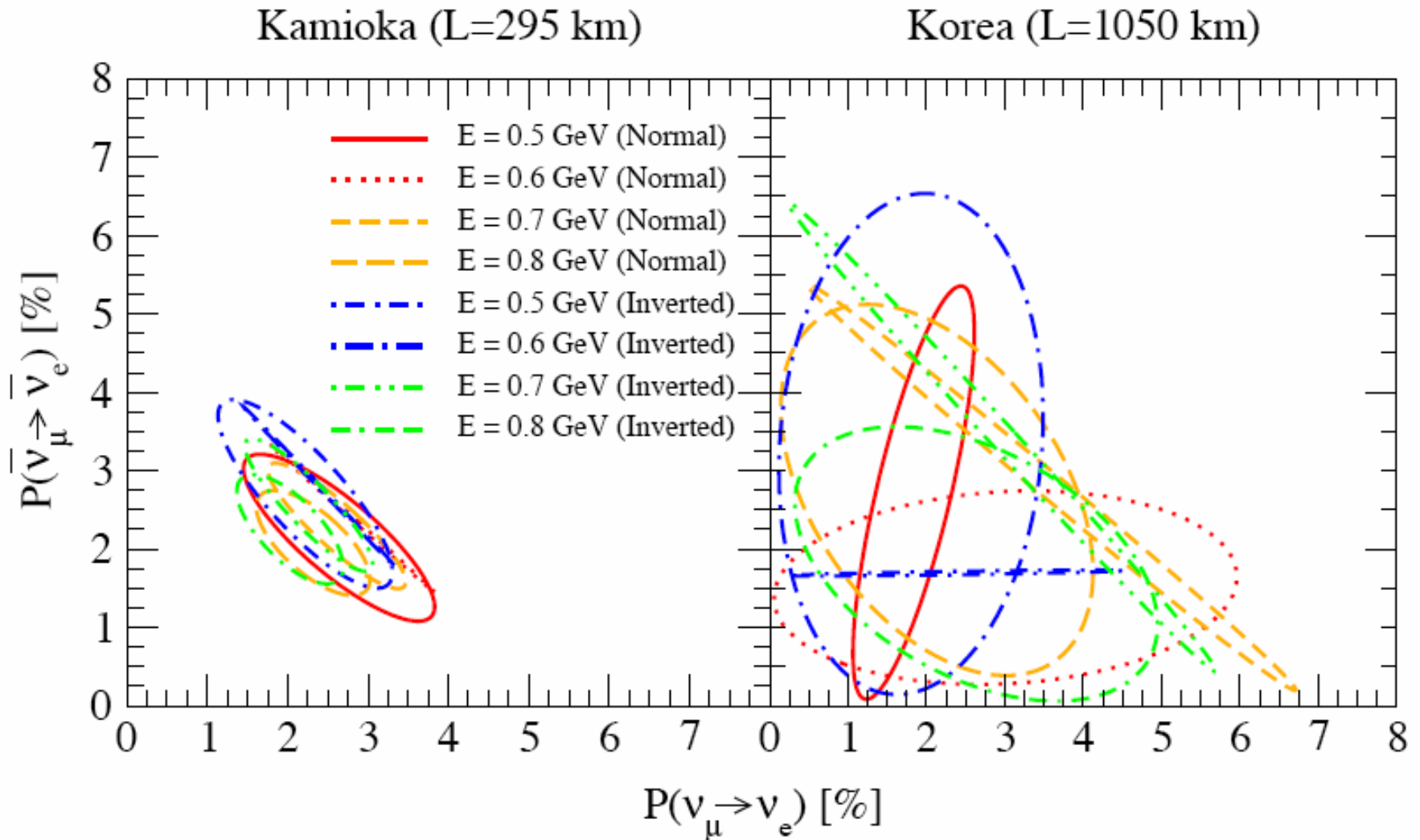
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# #1. Current design of Hyper-Kamiokande contains 2 tanks !

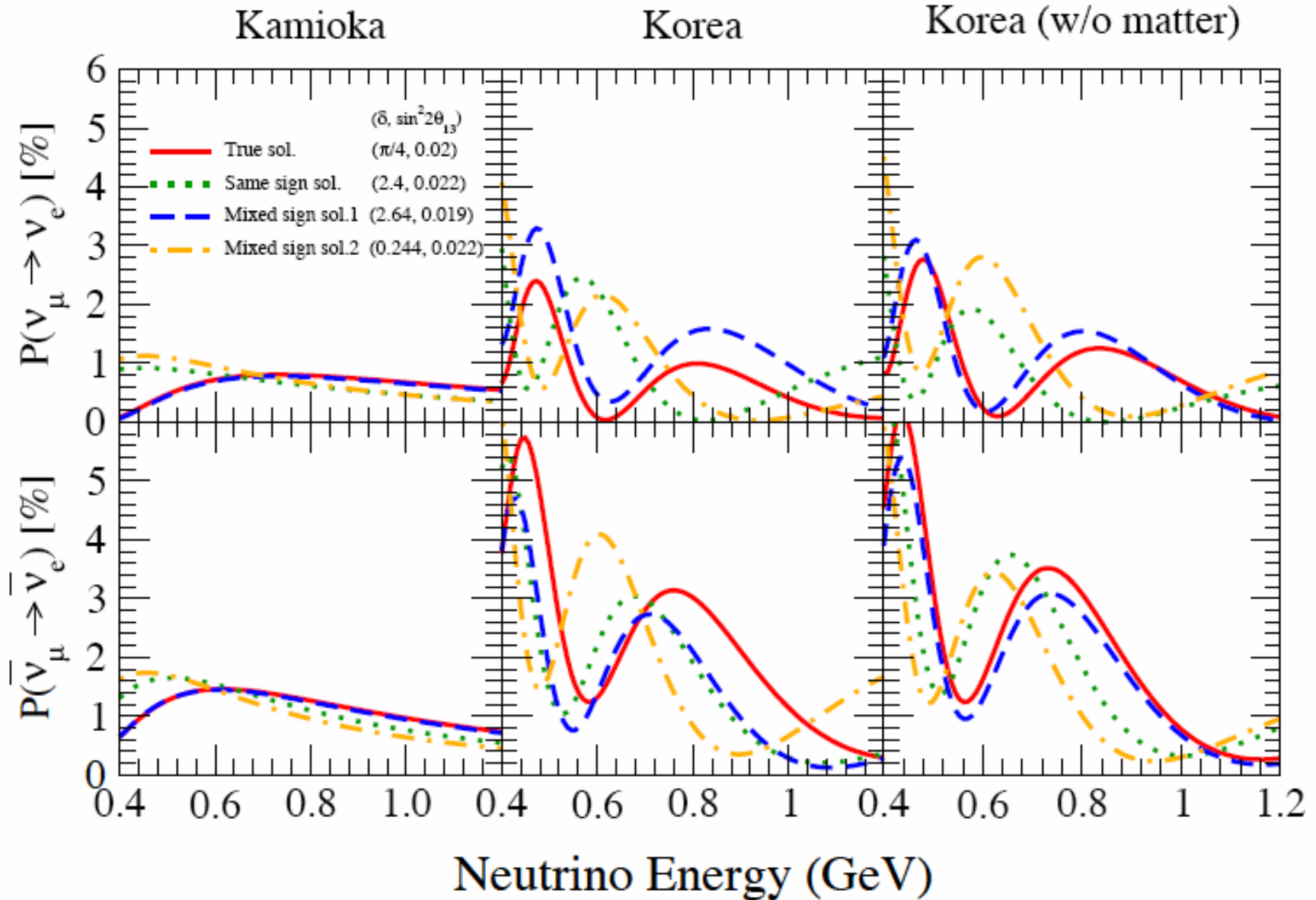


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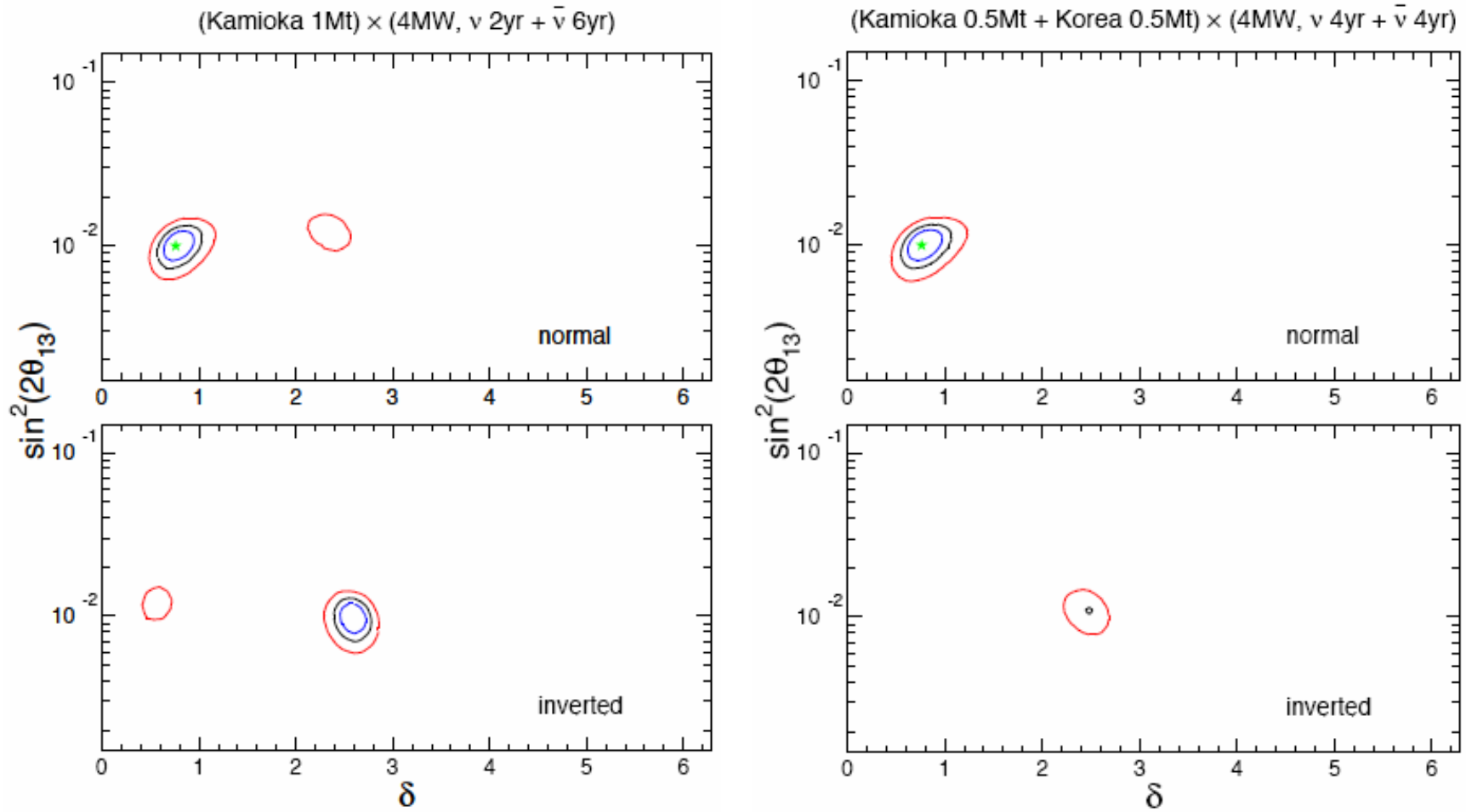
Sensitive to  $\delta$  because energy dependence  
is far more dynamic in 2nd oscillation  
maximum



# Spectral information solves degeneracy



# Spectral information solves intrinsic degeneracy





# $\chi^2$ definition

detector x beam  
combination

e-like bins

$\mu$ -like bins

systematic  
error term

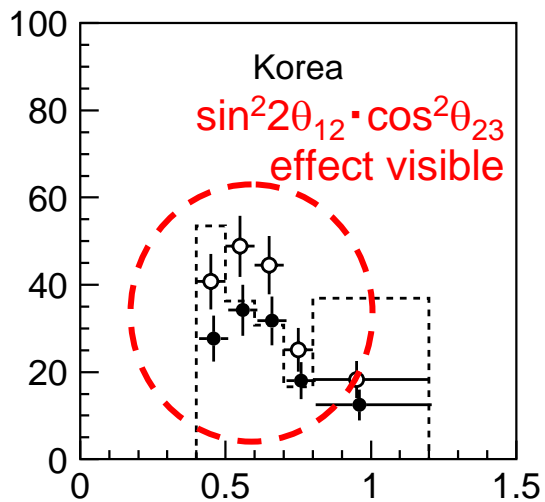
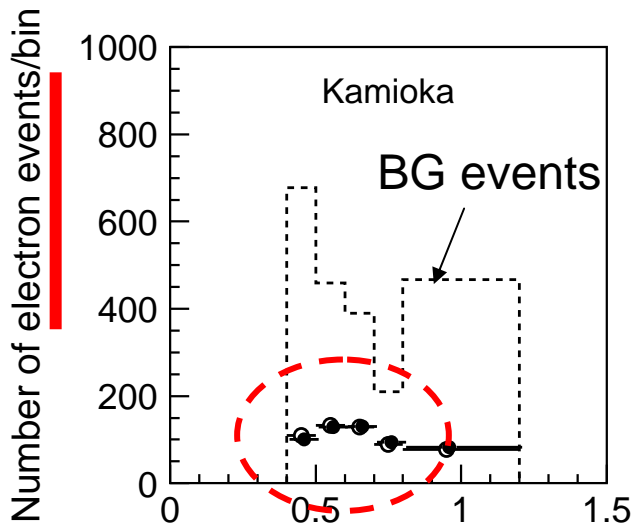
$$\chi^2 = \sum_{k=1}^4 \left( \sum_{i=1}^5 \frac{(N(e)_i^{\text{obs}} - N(e)_i^{\text{exp}})^2}{\sigma(e)_i^2} + \sum_{i=1}^{20} \frac{(N(\mu)_i^{\text{obs}} - N(\mu)_i^{\text{exp}})^2}{\sigma(\mu)_i^2} \right) + \sum_{j=1}^7 \left( \frac{\epsilon_j}{\tilde{\sigma}_j} \right)^2$$

$$N(e)_i^{\text{exp}} = N(e)_i^{\text{BG}} \cdot \left( 1 + \sum_{j=1,2,7} f(e)_j^i \cdot \epsilon_j \right) + N(e)_i^{\text{signal}} \cdot \left( 1 + \sum_{j=3,7} f(e)_j^i \cdot \epsilon_j \right)$$

$$N(\mu)_i^{\text{exp}} = N(\mu)_i^{\text{BG}} \cdot \left( 1 + \sum_{j=4,5,7} f(\mu)_j^i \cdot \epsilon_j \right) + N(\mu)_i^{\text{signal}} \cdot \left( 1 + \sum_{j=5,6,7} f(\mu)_j^i \cdot \epsilon_j \right)$$

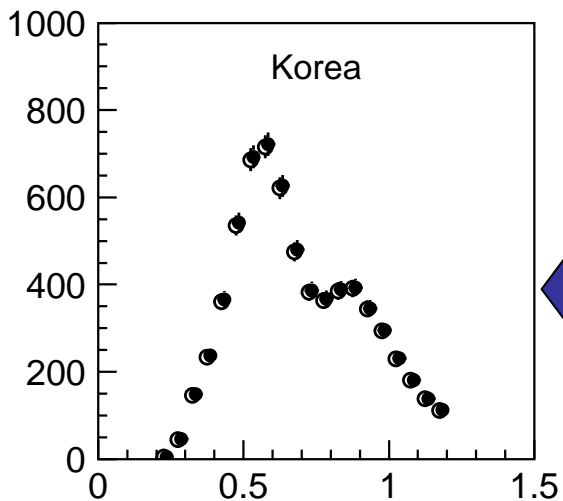
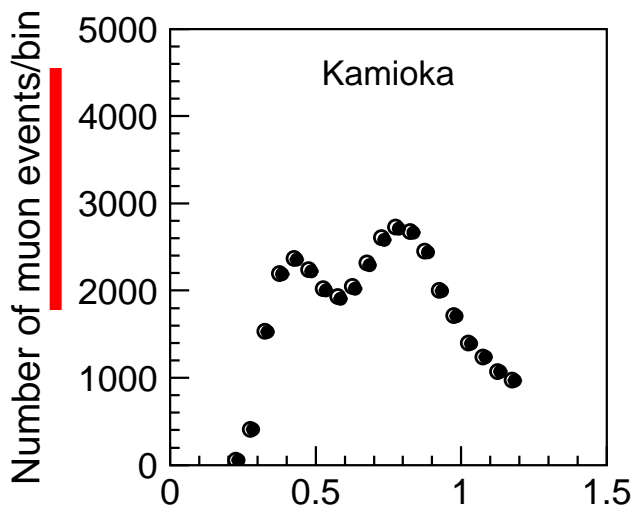
$f_j^i$  : fractional change in the predicted event rate in the  $i^{\text{th}}$  bin due to a variation of the parameter  $\epsilon_j$   
 $\epsilon_j$  : systematic error parameters, which are varied to minimize  $\chi^2$  for each choice of the oscillation parameters

# Effect of the solar term



○  $\sin^2 \theta_{23} = 0.4$ ,  
 $\sin^2 2\theta_{13} = 0.01$   
 ●  $\sin^2 \theta_{23} = 0.6$ ,  
 $\sin^2 2\theta_{13} = 0.0067$   
 (These parameters are chosen so that  $\sin^2 \theta_{23} \cdot \sin^2 2\theta_{13}$  is equal.)

$\Delta m^2$ : positive  
 $\delta = \pi 3/4$



← Included in this analysis, since  $\theta_{23}$  is relevant.

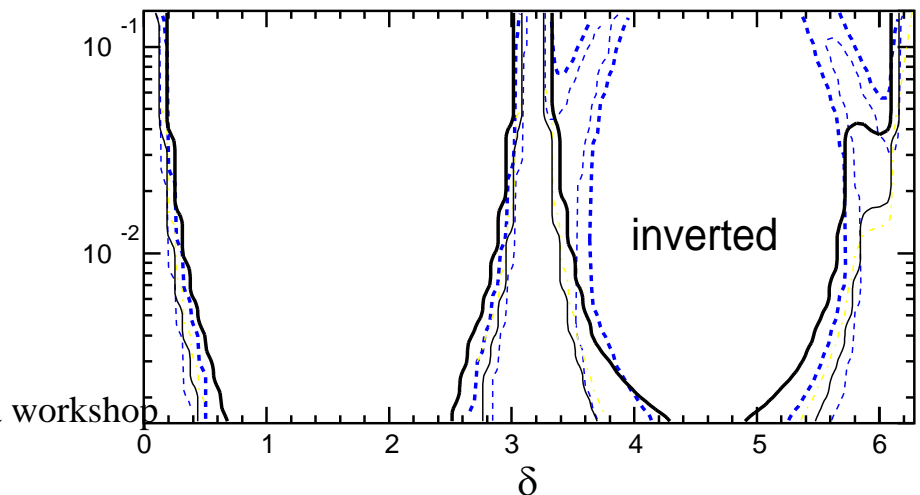
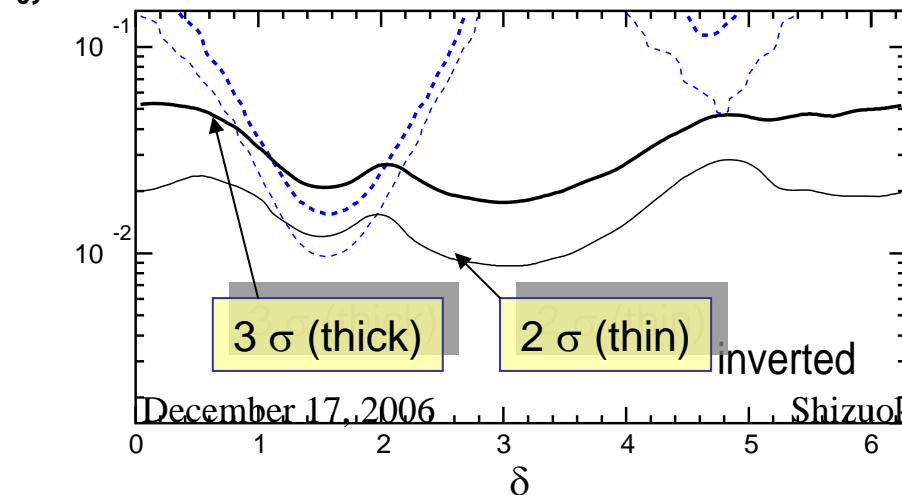
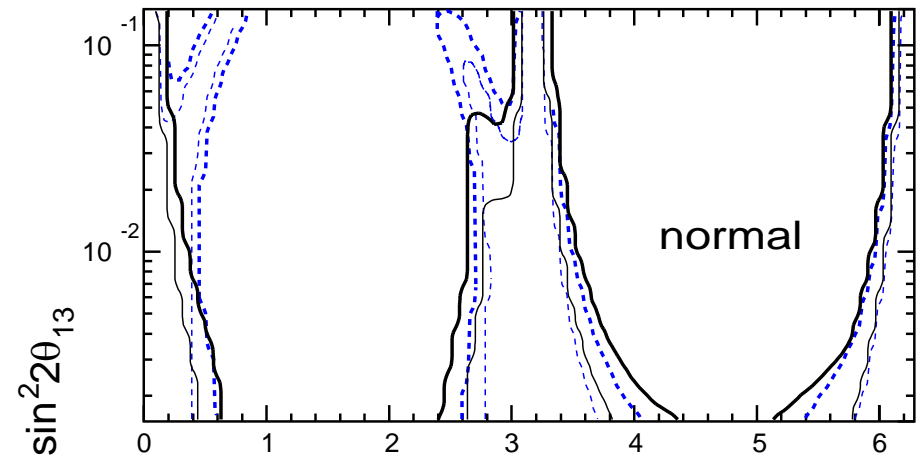
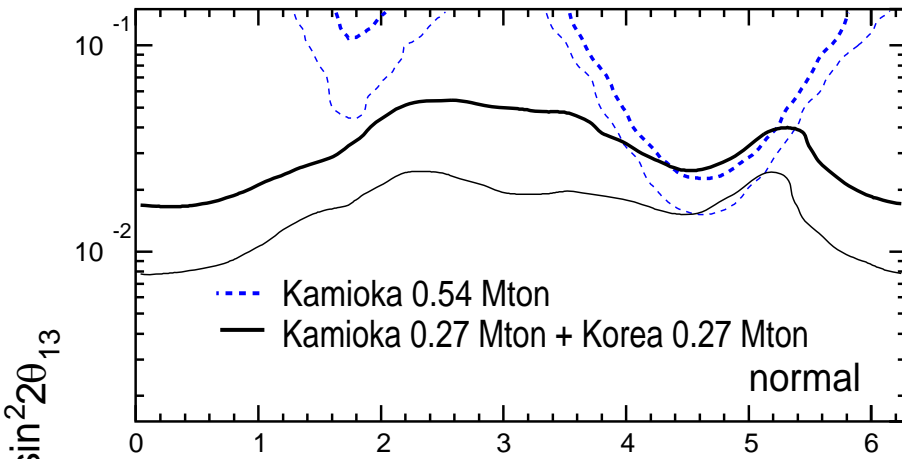
# T2KK vs. T2K II Comparison

hep-ph/0504026

Total mass of the detectors = 0.54 Mton fid. mass  
4 years neutrino beam + 4 years anti-neutrino beam

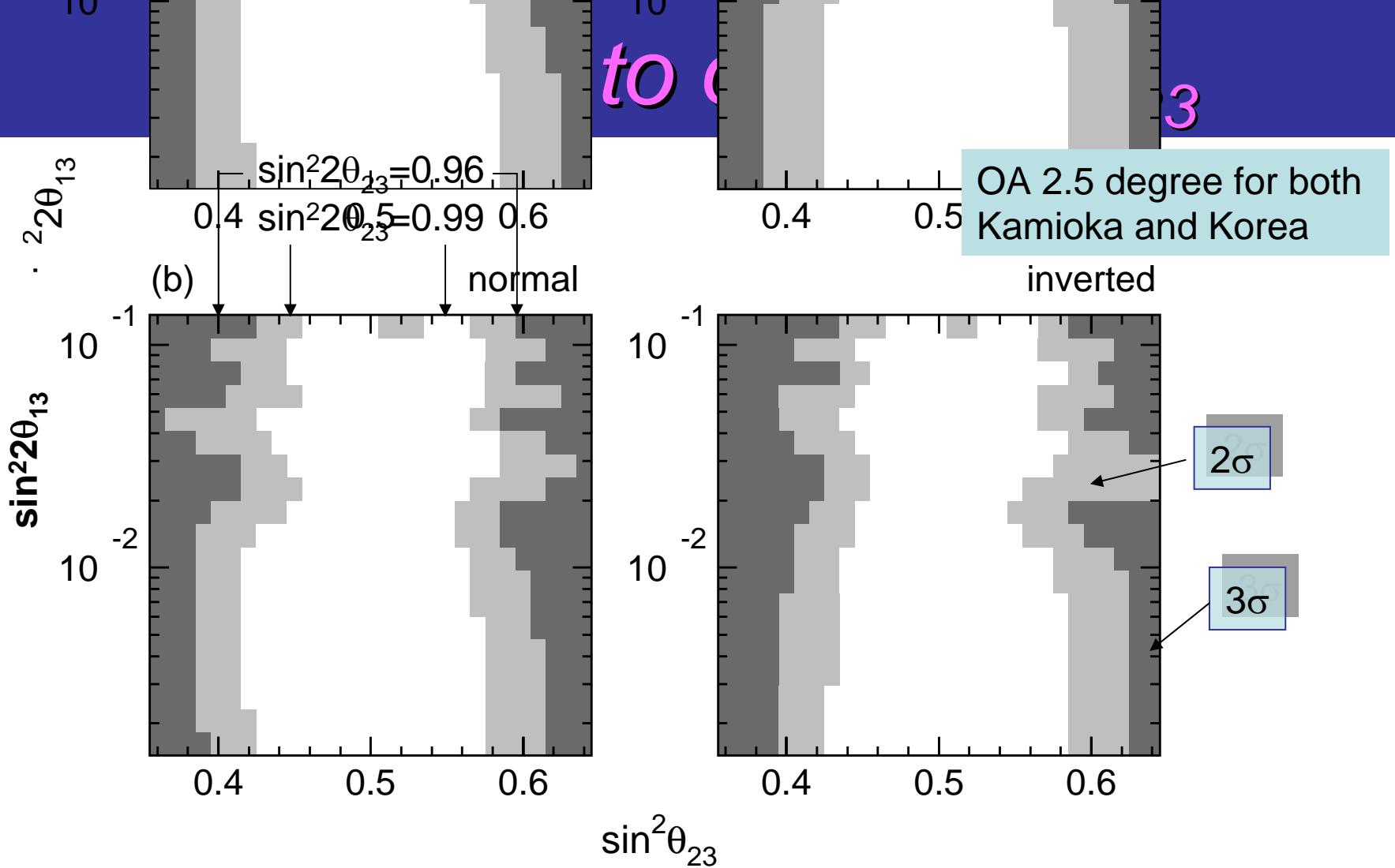
Mass hierarchy

CP violation ( $\sin\delta \neq 0$ )



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Octant ambiguity of  $\theta_{23}$  can be resolved if  $\sin^2 2\theta_{23} < \sim 0.97$  at  $2\sigma$  (almost independent of the value of  $\sin^2 2\theta_{13}$  and mass hierarchy).

➡ Can resolve the 8 fold degeneracy of the oscillation parameters.

# $\theta_{23}$ and sign- $\Delta m^2$ degeneracy decouple

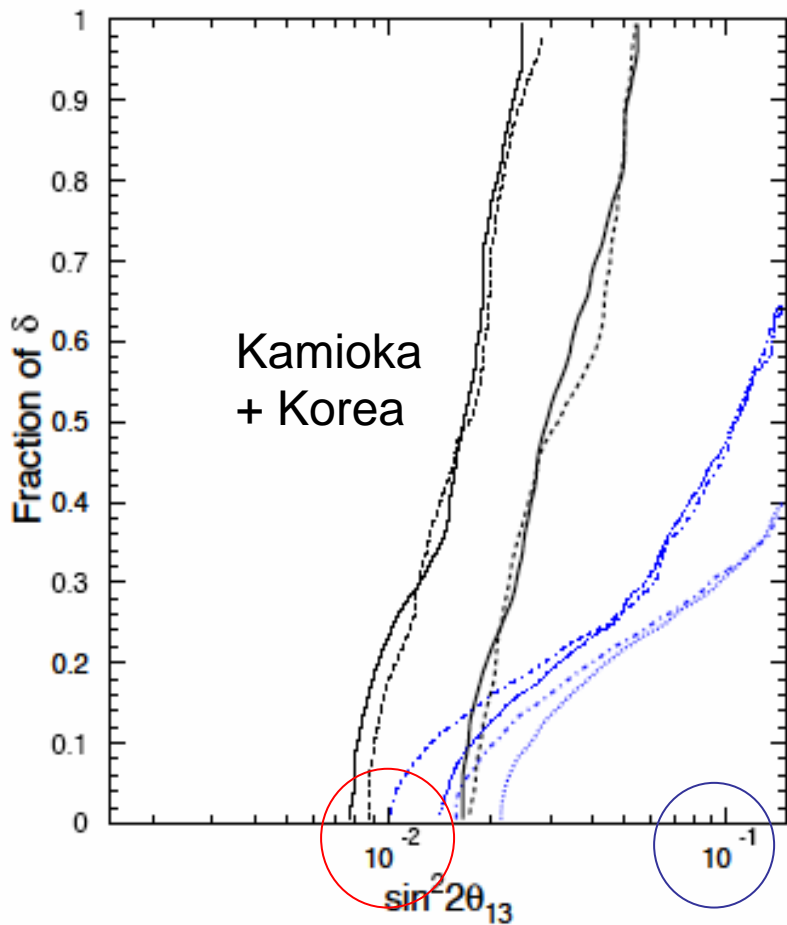
- For example, one can show, to first order in matter effect, the followings:
- $\Delta P(\text{octant}) = P(\text{1st octant}) - P(\text{2nd})$  is invariant under the interchange of two sign- $\Delta m^2$  degenerate pair
- $\Delta P(\text{hierarchy}) = P(\Delta m^2 +) - P(\Delta m^2 -)$  is invariant under the interchange of two  $\theta_{23}$  octant degenerate pair
- in T2K or T2KK setting, the intrinsic degeneracy is resolved by spectrum analysis decouple from the game 

In a nutshell, 8 fold degeneracy can be resolved by T2KK because ..

- intrinsic degeneracy is resolved by spectrum information
- sign- $\Delta m^2$  degeneracy is solved with matter effect + 2 identical detector comparison
- $\theta_{23}$  octant degeneracy is solved by identifying the solar oscillation effect in T2KK

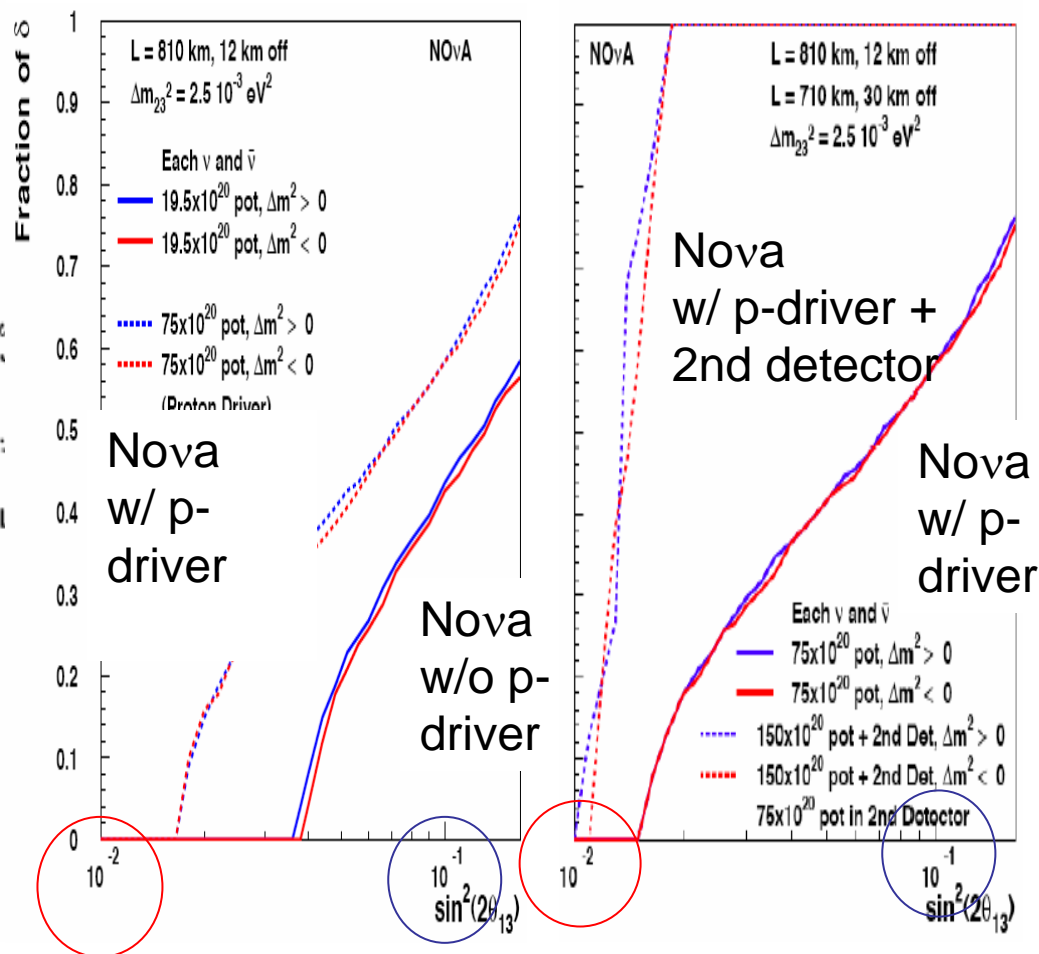
# Sensitivity to mass hierarchy: T2K-II vs. (Kam+Korea) vs. Nova

Sensitivity to mass hierarchy



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2  $\sigma$  Resolution of the Mass Hierarchy



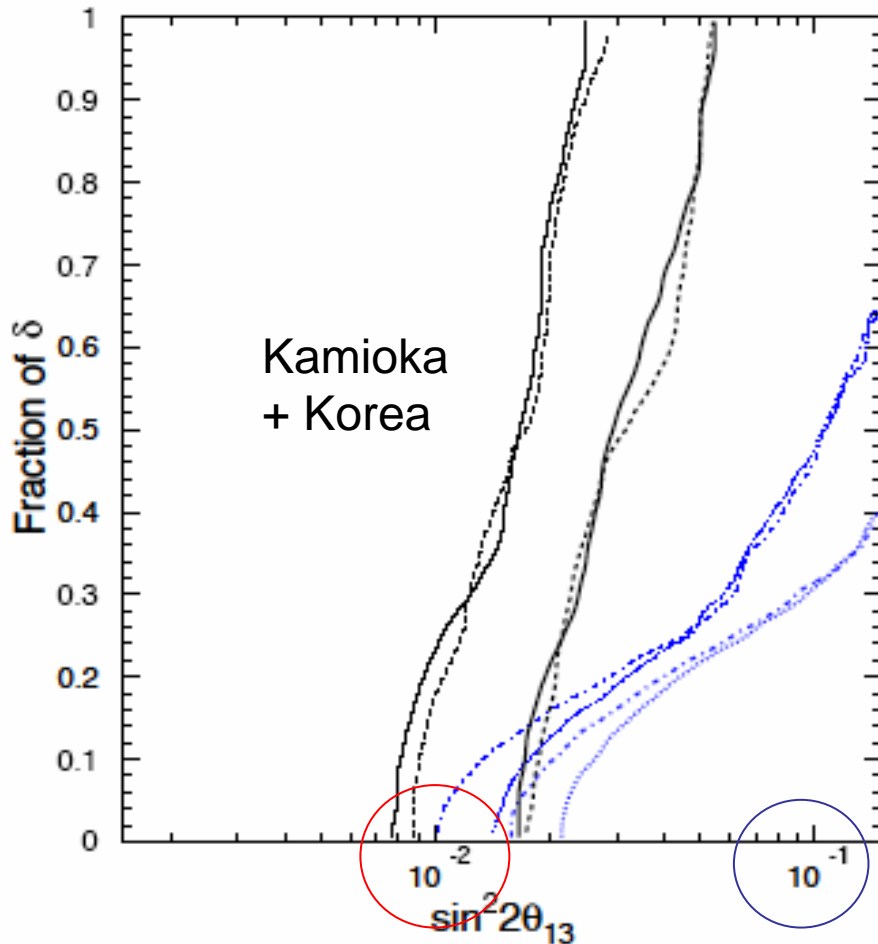
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# Expected sensitivity (2)

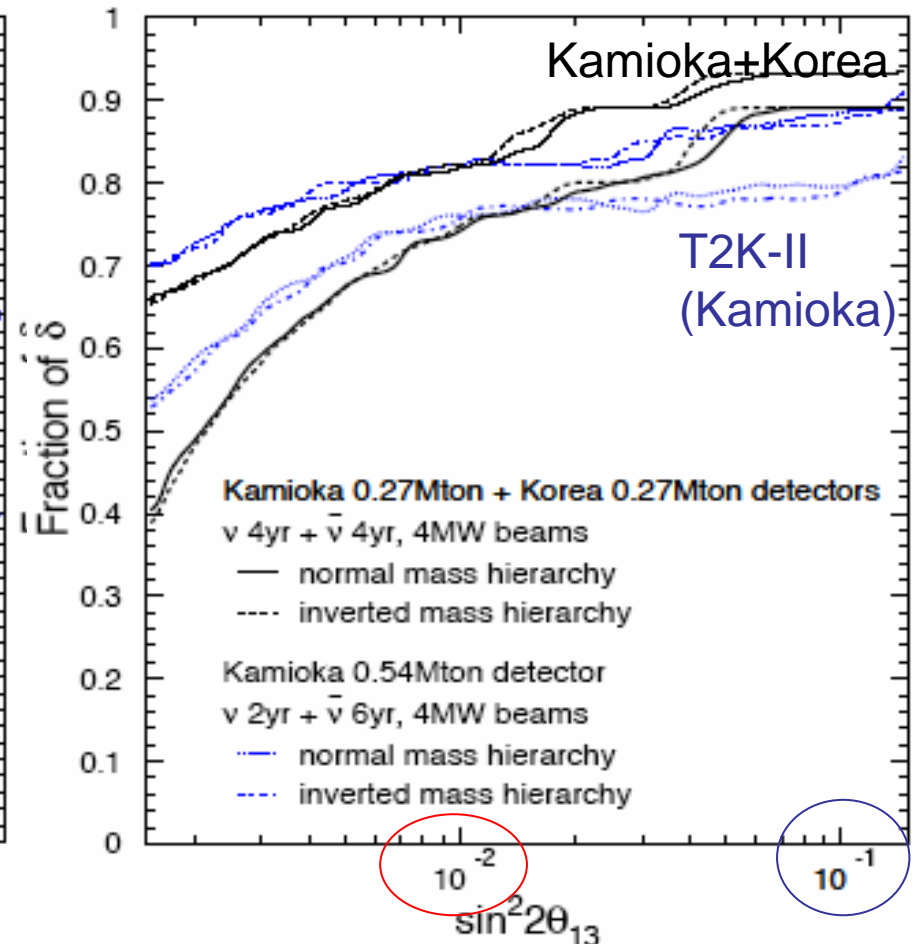
hep-ph/0504026

Total mass of the detectors = 0.54 Mton fid. mass  
4 years neutrino beam + 4 years anti-neutrino beam

Mass hierarchy




CP violation ( $\sin\delta \neq 0$ )





# T2KK; points of emphasis

- Observing Leptonic CP violation is a challenging goal  projects have to be expensive
- Sensitivity estimate has to be a reliable one, otherwise ...
- **Restrict to:** known background rejection technology by SK + conservative estimate of the systematic errors (5%) + identical 2 detector setting

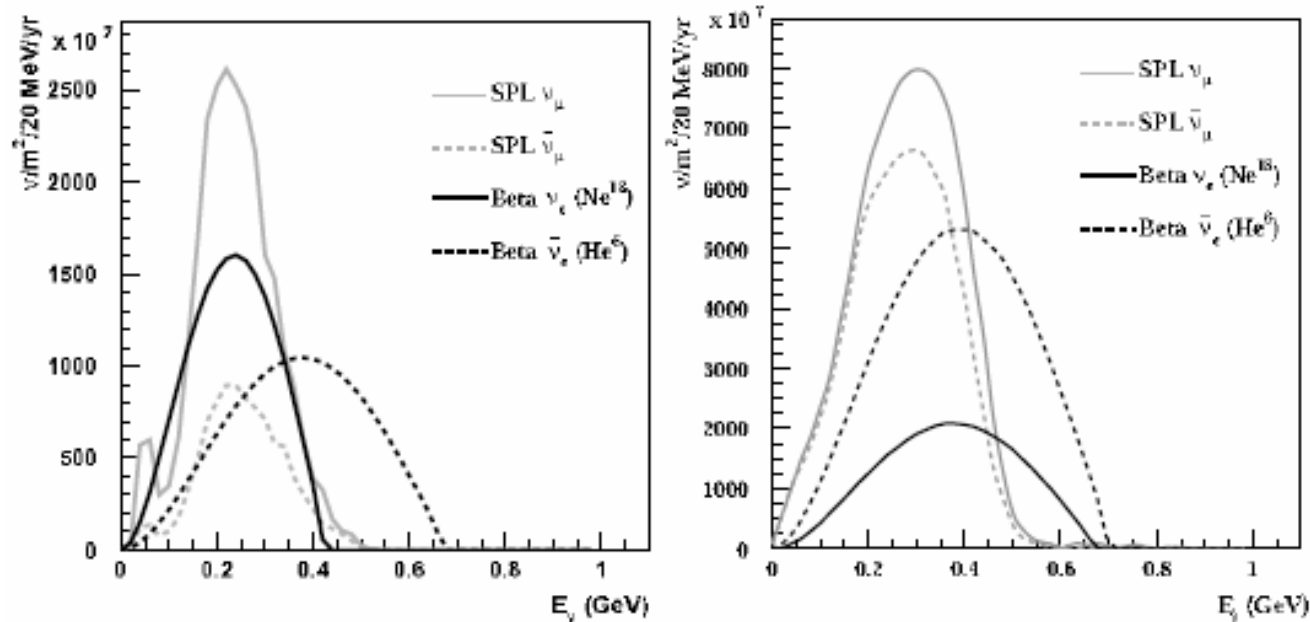
# Beta beam



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# Beta beam in a word



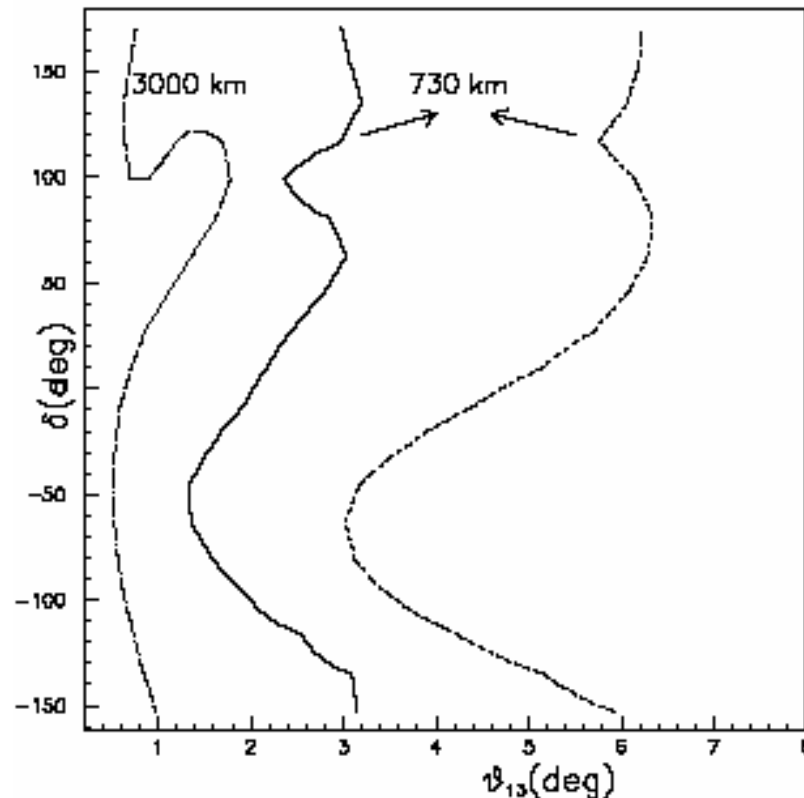
**Figure 6.** Comparison of neutrino fluxes from a super-beam (SPL) and a beta-beam. The neutrino beams are produced at CERN and sent to the Fréjus Underground Laboratory, 130 km from CERN. Two options for the beta-beam are shown here. Left: The ions circulate together in the storage ring, with  $\gamma = 60$  (100) for  $^6\text{He}$  ( $^{18}\text{Ne}$ ) (Mezzetto 2005). Right: The ions circulate at the same  $\gamma = 100$ , independently, in the storage ring. Note that the average neutrino energies are related to the ion boost through  $E_\nu \approx 2\gamma Q_\beta$  (Guglielmi *et al* 2005).

# What is good in Beta beam?

- pure  $\nu_e$  ( $^{18}\text{Ne}$ ) or  $\bar{\nu}_e$  ( $^6\text{He}$ ) beam
- charged pion background seems tolerable
- e- $\mu$  separation required but no charge ID required
- multi-MW proton beam NOT required

# Low vs. high $\gamma$ beta beam

- Setup I, low energy:  $\gamma = 60$  for  ${}^6\text{He}$  and  $\gamma = 100$  for  ${}^{18}\text{Ne}$ , with  $L = 130$  km (CERN–Fréjus) as in [12, 22].<sup>7</sup>
- Setup II, medium energy:  $\gamma = 350$  for  ${}^6\text{He}$  and  $\gamma = 580$  for  ${}^{18}\text{Ne}$ , with  $L = 732$  km (e.g. CERN–Gran Sasso with a refurbished SPS or with the LHC, FNAL–Soudan).
- Setup III, high energy:  $\gamma = 1500$  for  ${}^6\text{He}$  and  $\gamma = 2500$  for  ${}^{18}\text{Ne}$ , with  $L = 3000$  km (e.g. CERN–Canary islands with the LHC).



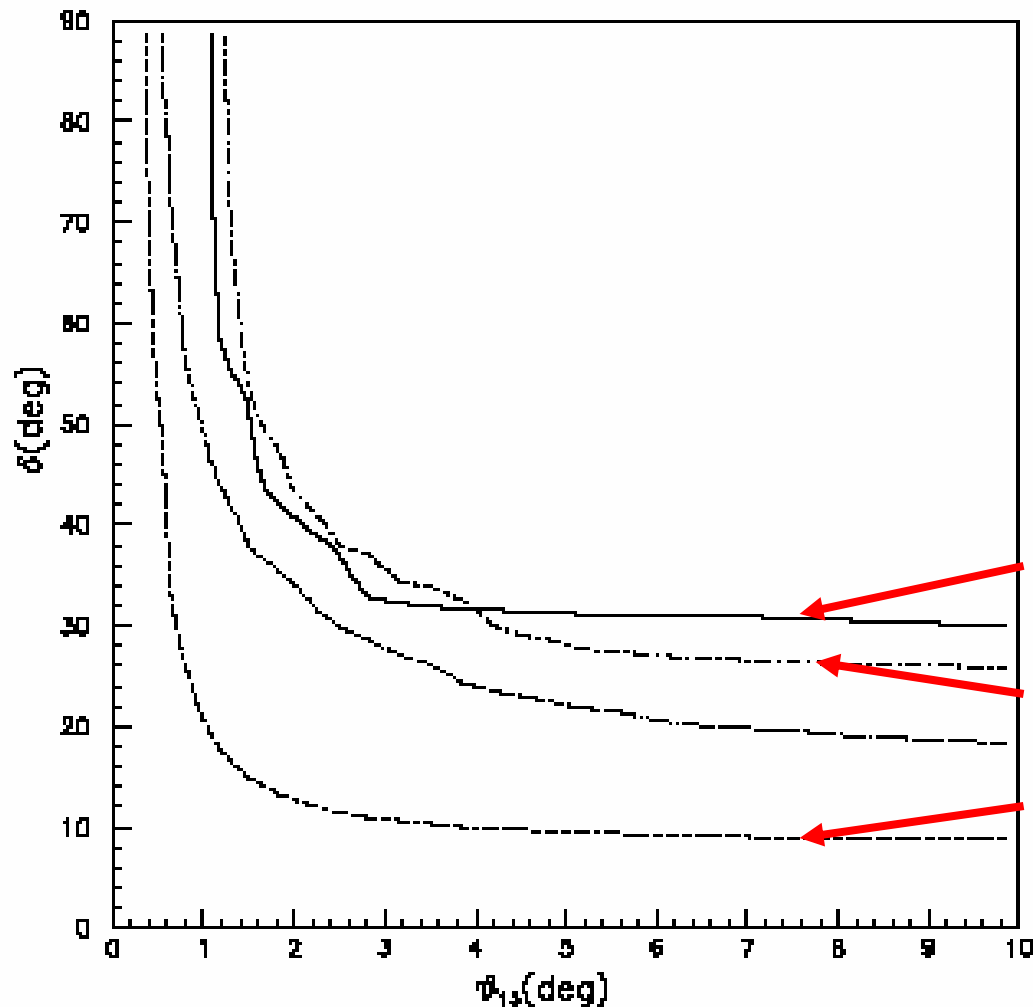
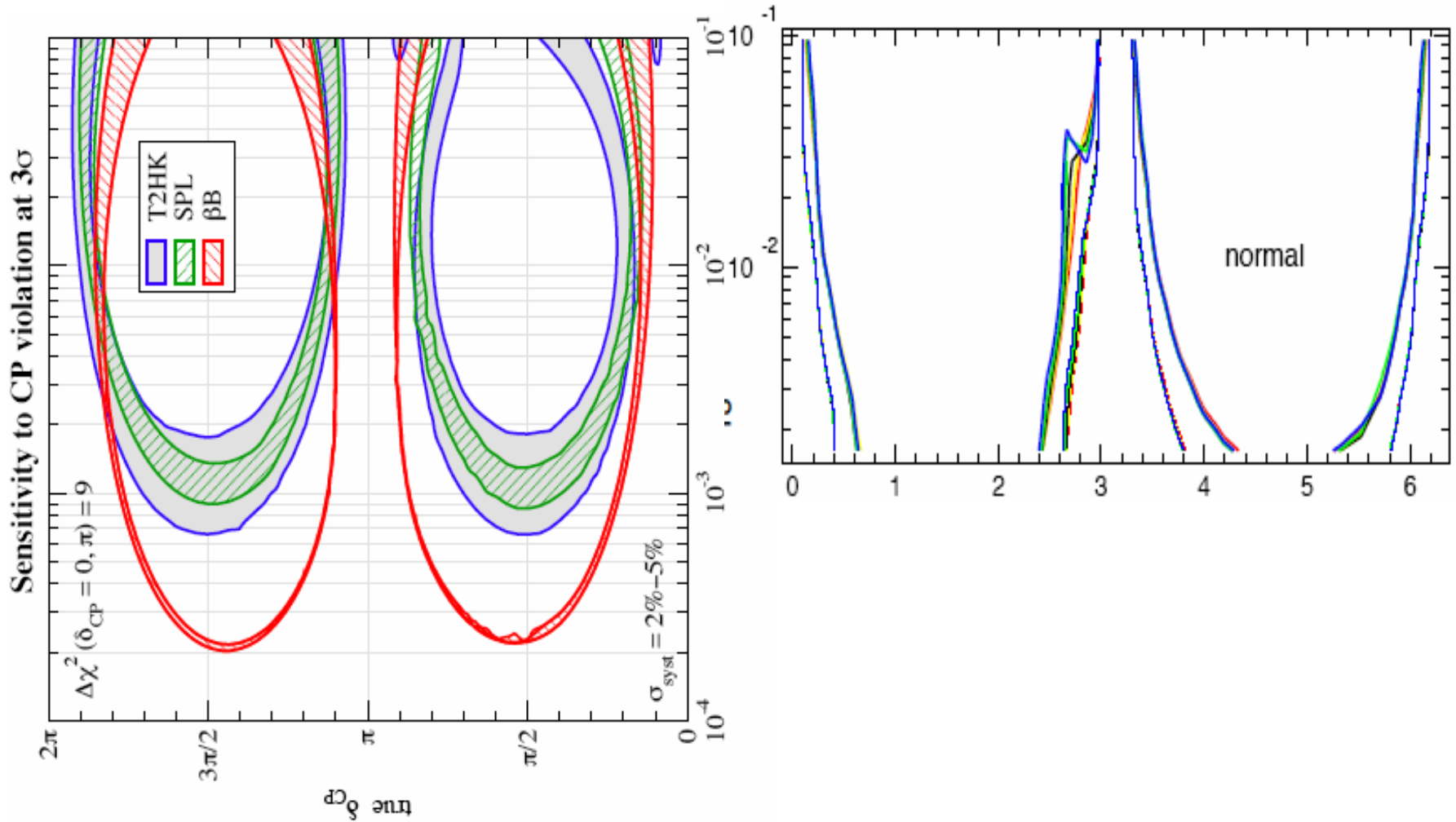


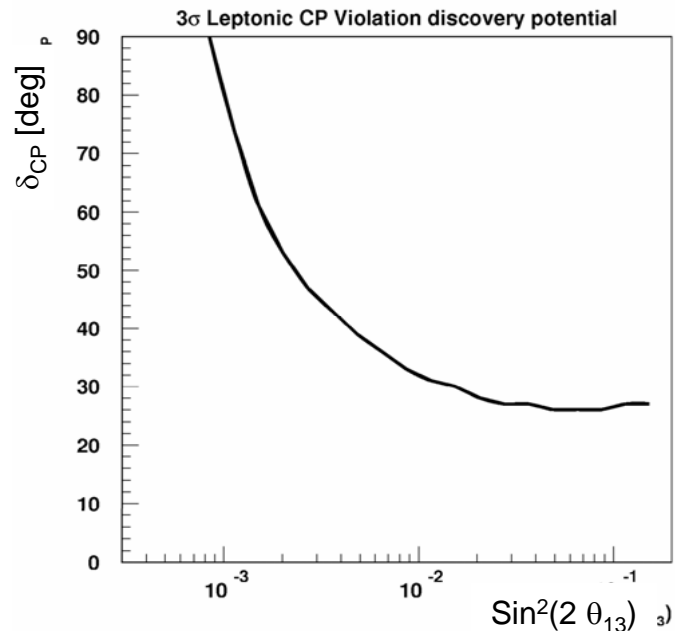
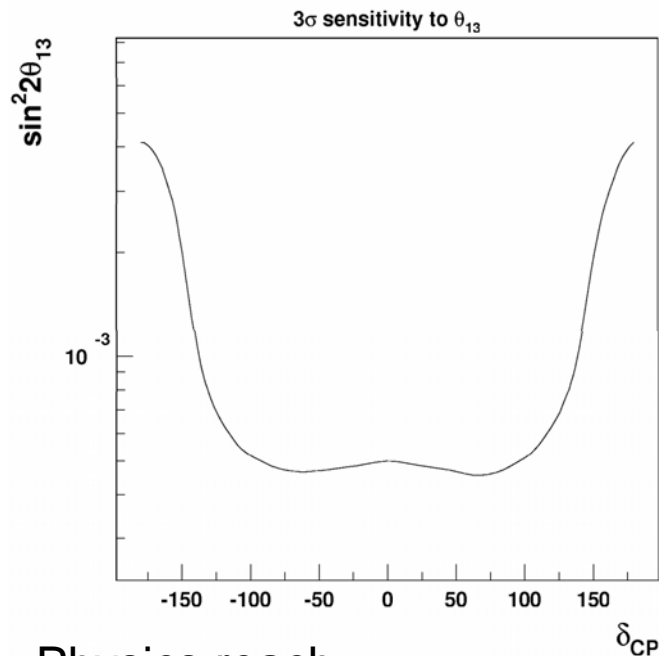
Figure 13: Region where  $\delta$  can be distinguished from  $\delta = 0$  or  $\delta = 180^\circ$  with a 99% CL for setup I (solid), setup II with the UNO-type detector of 400 kton described in section 3.1 (dashed) and with the same detector with a factor 10 smaller mass (dashed-dotted) and setup III (dotted) with a 40 kton tracking calorimeter described in section 3.4.

# Beta vs. T2KK



# Physics reach; low $\gamma$

- EURISOL scenario
  - $\gamma=100$
  - each  ${}^6\text{He}$  and  ${}^{18}\text{Ne}$  with a 5-year run
  - $2.9 \cdot 10^{18}$   ${}^6\text{He}$  decays/year or  $1.1 \cdot 10^{18}$   ${}^6\text{Ne}$  decays/year



- Physics reach
  - Sensitivity on  $\Theta_{13}$  down to  $\sim 1^\circ$



# Neutrino factory

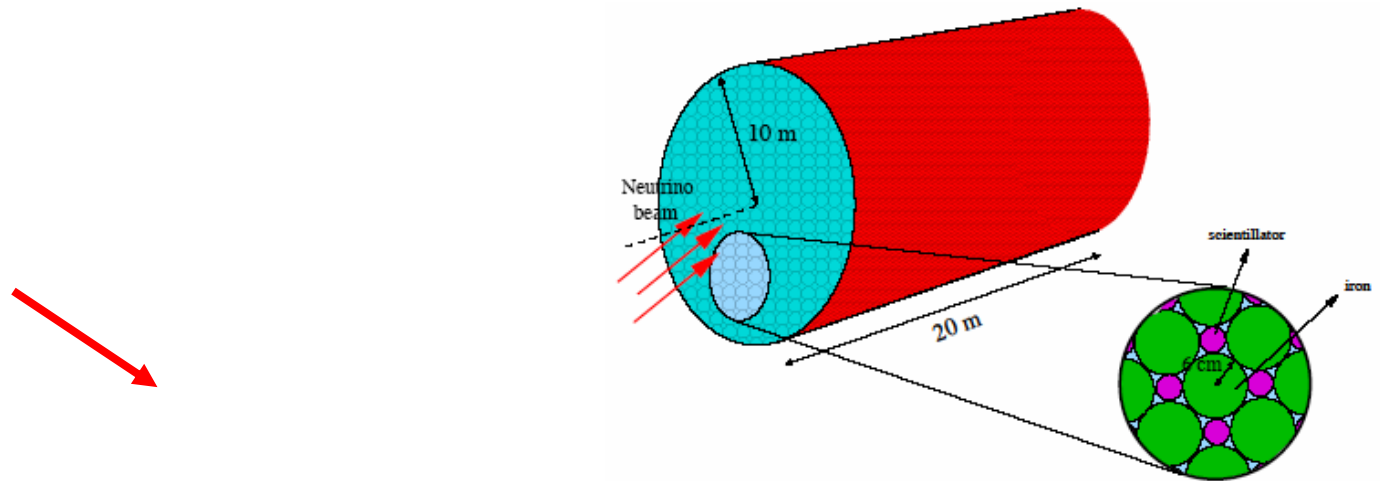


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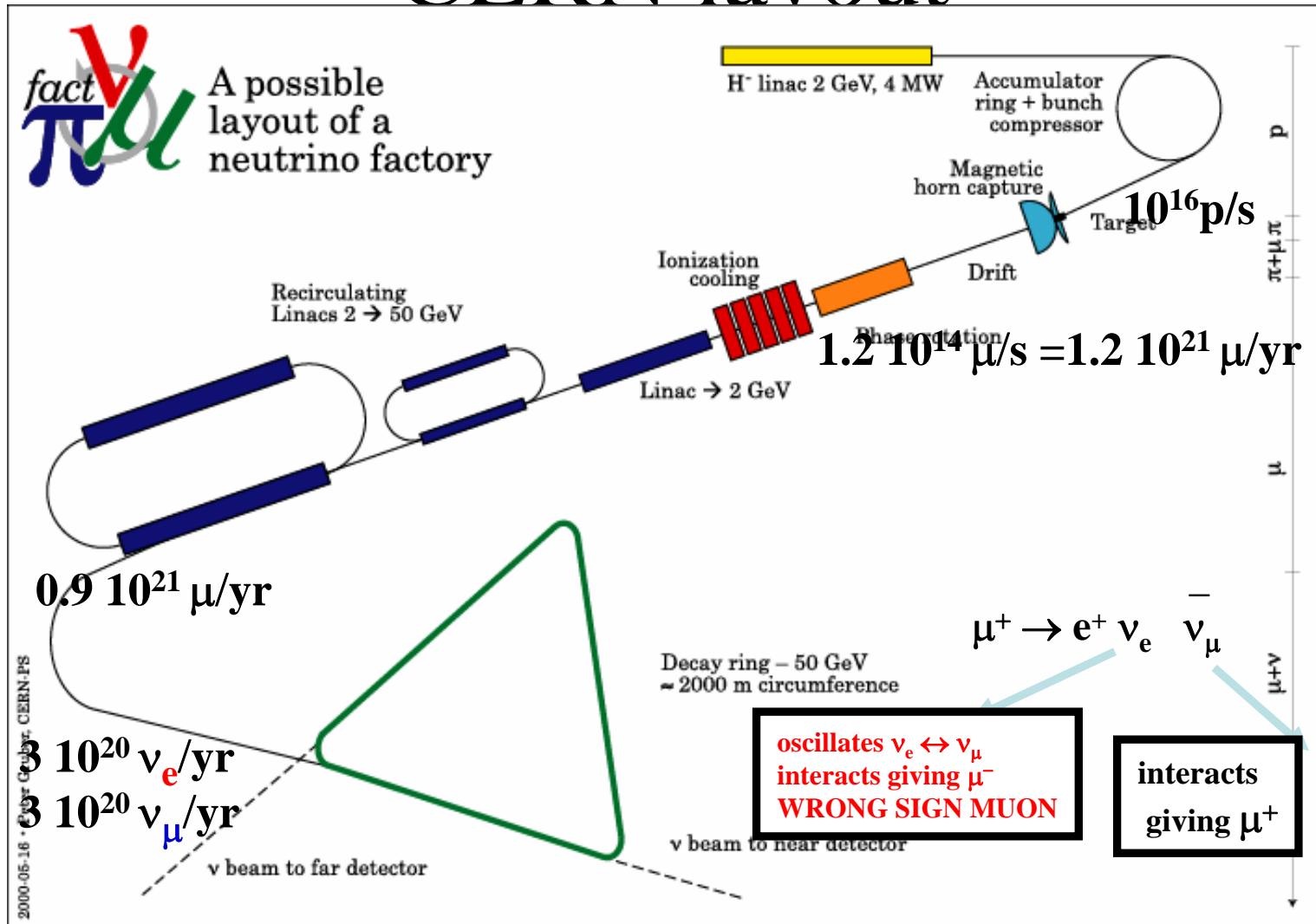
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# What is good in Neutrino factory ?

- well understood combination of  $\nu_e$  and  $\nu_\mu$  beam with precisely ( $\sim 10^{-5}$ ) known muon energy
- small background (how small,  $10^{-4} - 10^{-5}$  ?)
- muon charge ID required
- multi-MW proton beam required



# -- Neutrino Factory -- CERN layout



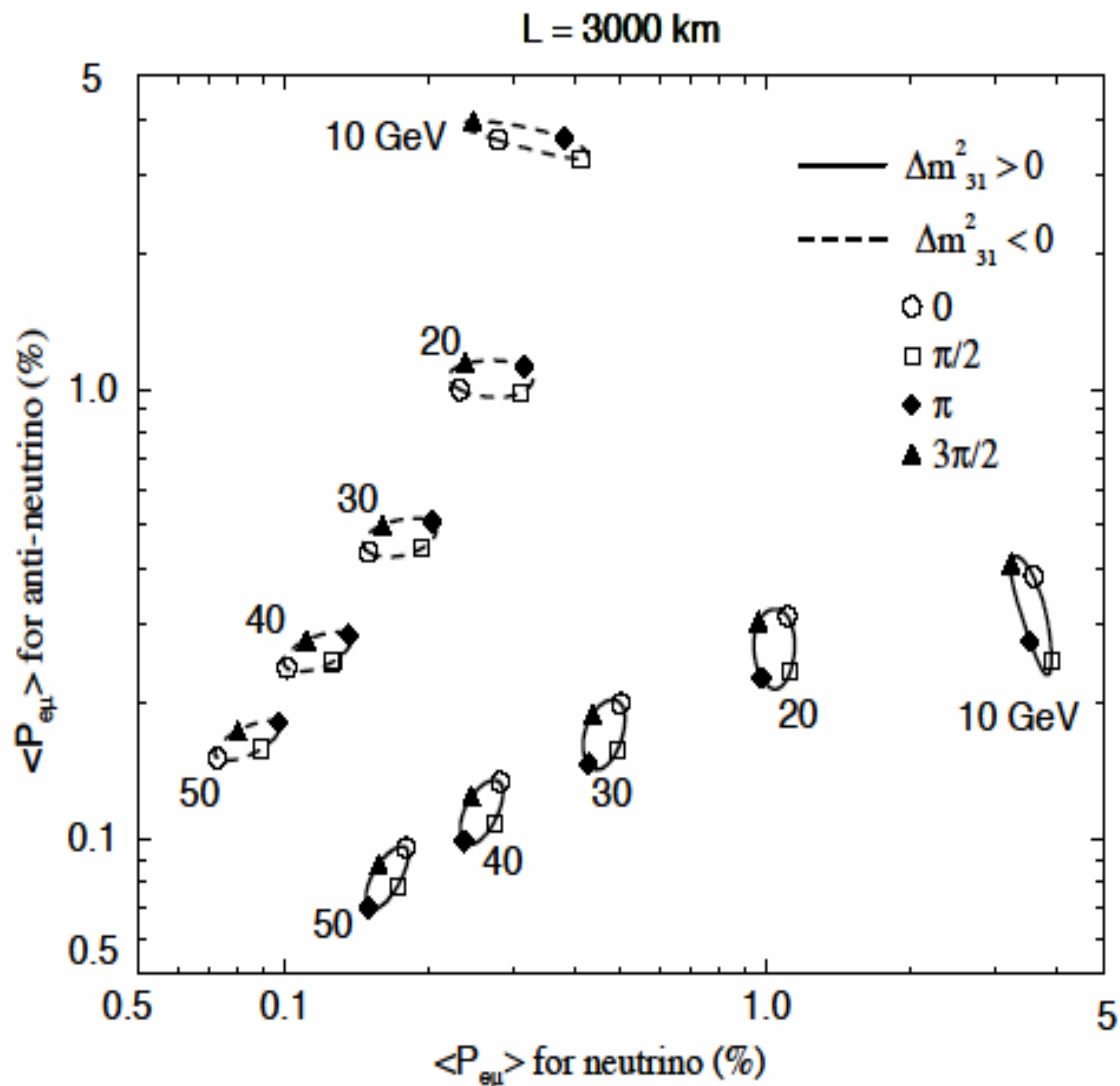
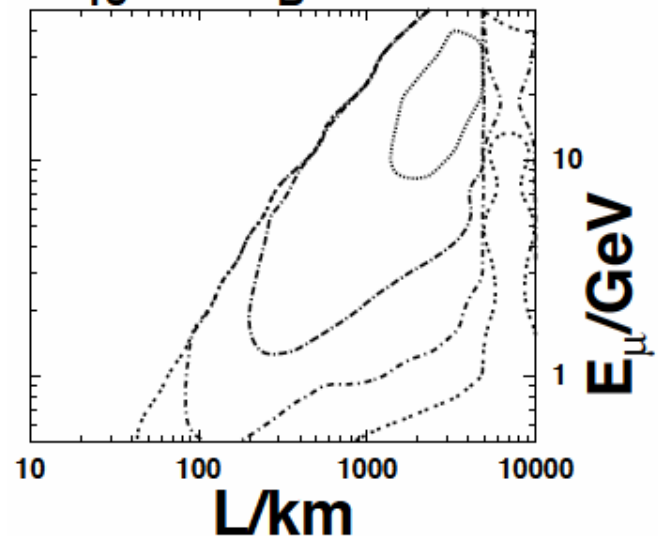


Figure 8: The CP trajectory diagram in bi-probability plane for  $L = 3000 \text{ km}$  and much higher neutrino energies  $E = 10 - 50 \text{ GeV}$  which correspond to so called “Neutrino Factory” situation. The mixing parameters are fixed to be the same as in figure 1 except that we take  $\rho Y_e = 2.0 \text{ g/cm}^3$ .

# Optimal energy E & baseline L

- $E = 30 \sim 50 \text{ GeV}$
- $L \sim 3000 \text{ km}$
- # of events =  $(E^2 / L^2) \times E \times (L/E)$

$$= (E^2 / L^2) \times L = (E / L) \times E \quad \theta_{13}=1^\circ \quad f_B=10^{-3}$$



# Magic baseline or refraction length

$$P(\nu_\mu \rightarrow \nu_e) = |\sqrt{P_{atm}} + e^{i(\delta \pm \frac{\Delta_{31}}{2})} \sqrt{P_{solar}}|^2$$

At  $aL=2\pi$ ,  $P_{solar}$   
vanishes -->  
No CP phase  
dependence



$$P_{atm} = \left( s_{13} s_{23} \Delta_{31} \frac{\sin\left(\frac{\Delta_{31} \mp aL}{2}\right)}{\left(\frac{\Delta_{31} \mp aL}{2}\right)} \right)^2$$

$$P_{solar} = \left( c_{12} s_{12} c_{23} \Delta_{21} \frac{\sin\left(\frac{aL}{2}\right)}{\left(\frac{aL}{2}\right)} \right)^2$$

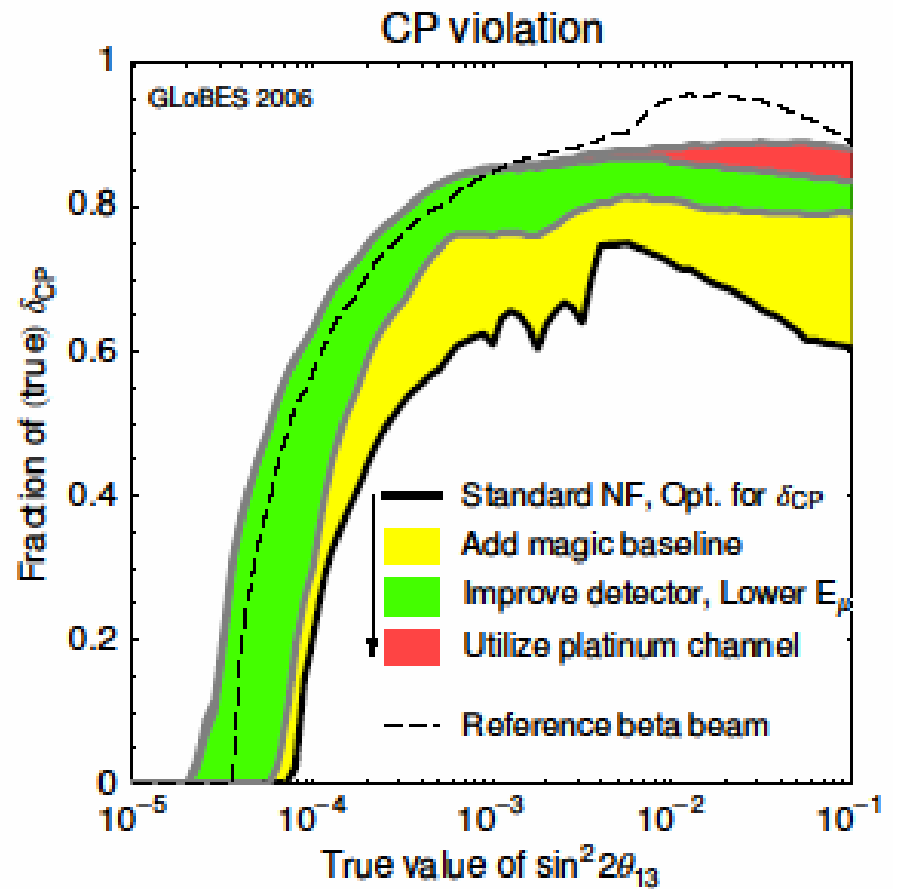
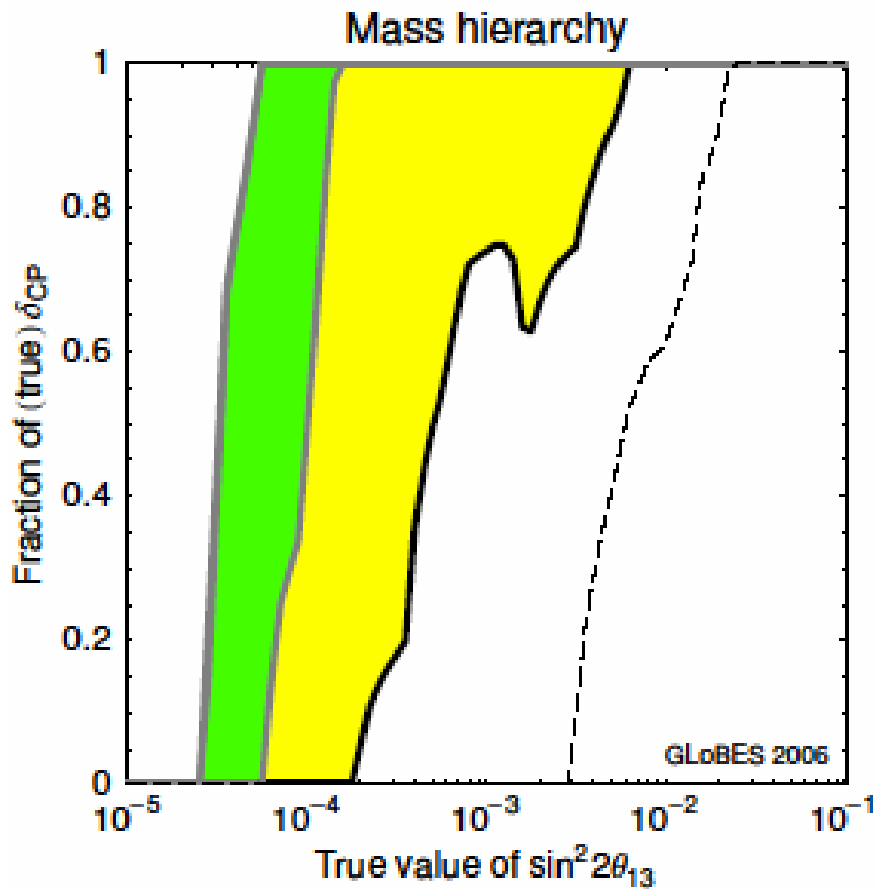
$$\Delta_{31} \equiv \frac{|\Delta m_{31}^2| L}{2E}, \quad a = \sqrt{2} G_F N_e(x),$$

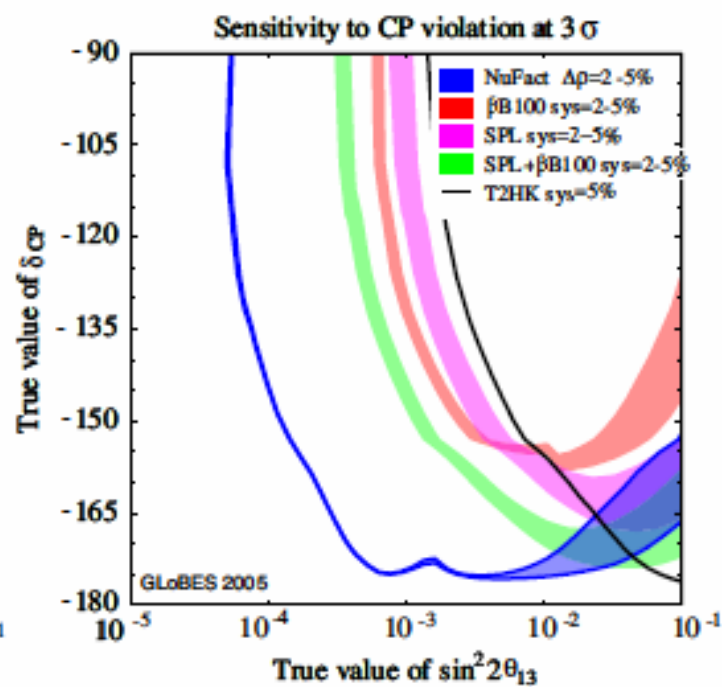
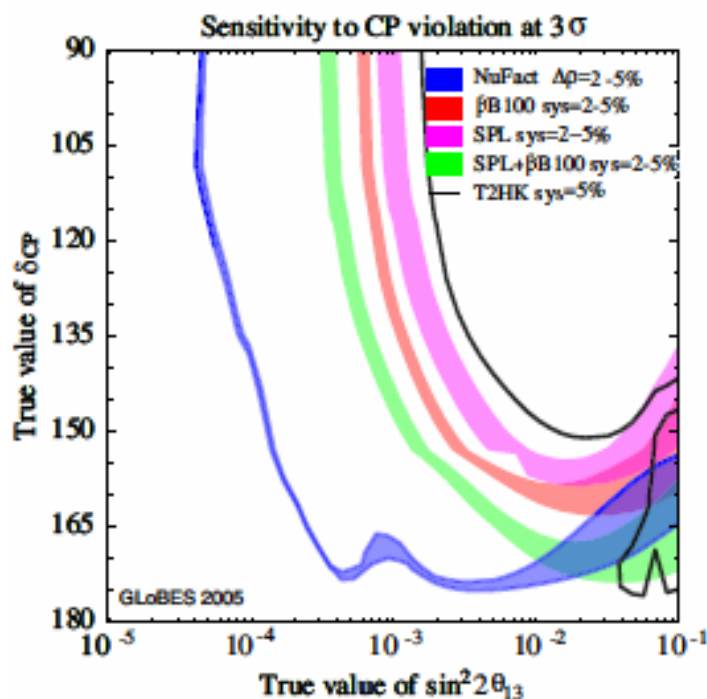
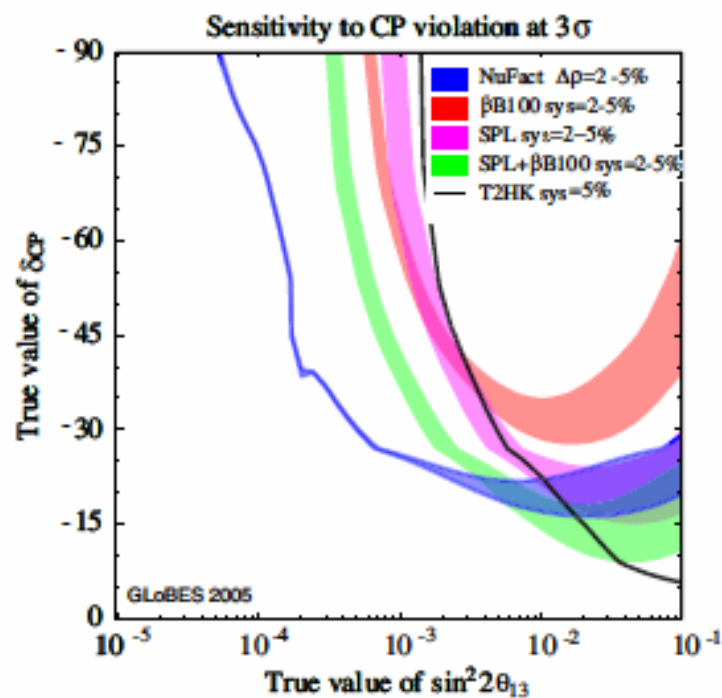
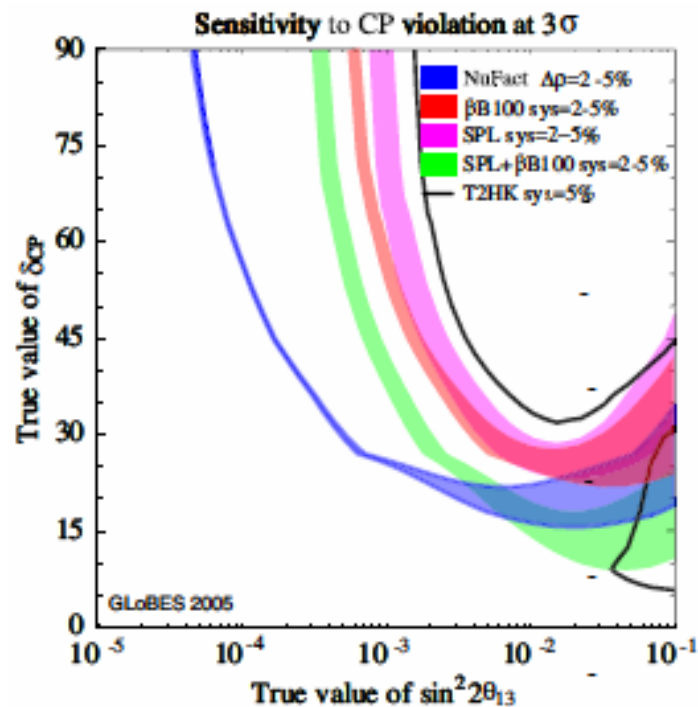
$$\pm = \text{sign of } \Delta m_{31}^2$$

“magic baseline”

Help resolve  
degeneracy

# Nufact sensitivity







# Conclusion

- The next-generation and some future options for LBL experiments are reviewed
- still long way to complete the MNS matrix;  $\theta_{13}$  first, and then  $\delta$  and mass hierarchy
- T2KK is powerful enough to solve 8-fold parameter degeneracy in situ
- if  $\theta_{13} < 3^\circ$ , we need  $\beta$  beam and/or neutrino factory; the choice is highly debatable -> exciting possibilities because the small  $\theta_{13}$  may imply "symmetry"

# Koide-san and teQuila



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